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# Luminosity and stellar mass functions of disks and spheroids in the SDSS and the supermassive black hole mass function

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## ABSTRACT

Using the GALACTICA code of Benson et al., we obtain quantitative measurements of spheroid-to-disk ratios for a sample of 8839 galaxies observed in the Sloan Digital Sky Survey. We carry out extensive tests of this code and of GIM2D, finding that they perform similarly in all respects. From the spheroid and disk luminosities, we construct luminosity and stellar mass functions for each component and estimate the relative luminosity and stellar mass densities of disks and spheroids in the local Universe. Assuming a simple one-to-one mapping between spheroid mass and the mass of a central supermassive black hole, we provide the most accurate determination so far of the black hole mass function in the local universe. From this, we infer a cosmological mass density of black holes of  $\rho_{\bullet} = (2.40 \pm 0.13) \times 10^5 h M_{\odot} \text{Mpc}^{-3}$ . We compare our results to predictions from current hierarchical models of galaxy formation and these are found to fare well in predicting the qualitative trends observed. In units of the critical density of the Universe, we find  $\Omega_{\text{stars,disks}} = (0.656 \pm 0.005) h^{-1} 10^{-3}$  and  $\Omega_{\text{stars,spheroids}} = (0.383 \pm 0.005) h^{-1} 10^{-3}$ , implying that the Universe contains almost twice as much stellar mass in disks as in the spheroidal components of galaxies.

**Key words:** galaxies: structure; galaxies: abundances; galaxies: bulges; galaxies: luminosity function, mass function; galaxies: statistics

## 1 INTRODUCTION

The distinction between disks and spheroids is one of the defining properties of galaxies. Determining the relative importance of these two basic types of galactic component is fundamental to a broad characterisation of the galaxy population. Yet, this is a complicated task which requires not only high quality imaging for large samples, but also software capable of decomposing the light from each object into a disk and a spheroid<sup>1</sup>. The determination of spheroid luminosities has recently received even more prominence since the discovery that perhaps all galaxies harbour a supermassive black hole at their centre whose mass is proportional to the luminosity of the spheroid or bulge (Kormendy & Richstone 1995; Magorrian et al. 1998; Merritt & Ferrarese 2001; Marconi & Hunt 2003; Häring & Rix 2004).

From a theoretical point of view, explaining why most of the stars in the Universe end up either in disks or in

spheroids and understanding the physical processes that result in the formation of one or the other of these morphological structures is a major challenge. The current theoretical framework used to investigate galaxy formation is the cold dark matter model (Peebles 1982, Davis et al 1985), in which galaxies build up hierarchically. Within this model, the basic processes thought to be responsible for the distinction between disks and spheroids were identified over twenty years ago (Blumenthal et al. 1984; Frenk et al 1985): disks result from the collapse of rotating gas cooling within dark matter halos whereas spheroids result from major mergers or disk instabilities (Barnes & Hernquist 1992; Mo, Mao & White 1998). The traditional categories of galaxy morphology, spirals, irregulars, etc, are too detailed for current theoretical models to explain, but the relative luminosities and stellar masses of spheroids and disks can readily be predicted (Baugh et al. 1996a; Baugh et al. 1996b; Kauffmann, Charlot & White 1996; Kauffmann & Charlot 1998; Somerville & Primack 1999; Hatton et al. 2003). Thus, accurate measurements of these quantities, for example, as a function of

<sup>1</sup> We will use the term “spheroid” throughout to refer to both elliptical galaxies and the bulges of spiral galaxies

absolute magnitude and in different environments provides a powerful test of models of galaxy formation and evolution.

An early attempt to determine the relative contributions of spheroids and disks to the luminosity density of the Universe was made by Schechter & Dressler (1987). They studied a magnitude limited sample of  $\sim 200$  galaxies brighter than  $V = 16.5$ , drawn from the catalogue of Dressler (1980a), and determined spheroid-to-disk ratios by visual inspection. From this, they derived the distribution of spheroid-to-disk ratios, as a function of absolute magnitude, and found the overall spheroid-to-disk ratio to be higher in high density environments (galaxy clusters) than in low density environments (the “field”). Schechter & Dressler (1987) found that disks appear to contribute roughly twice as much as spheroids to the mean luminosity density of the Universe. Since a large fraction of the disk light comes from a relatively small number of young stars, Schechter & Dressler (1987) concluded that the relative contribution of spheroid and disk components to the mean stellar mass density of the Universe is very nearly equal.

More recently, Benson et al. (2002) developed a quantitative method to determine galaxy morphology, specifically to estimate spheroid-to-total (S/T) light ratios. This method is implemented in the code GALACTICA (GALaxy Automated Component Image Construction Algorithm). Benson et al. (2002) analysed a magnitude-limited sample of  $\sim 100$  field galaxies brighter than  $I = 16.0$  and found the luminosity functions of spheroids and disks to be remarkably similar. They provisionally concluded that spheroids and disks contribute almost equally to the total stellar mass density in the Universe but stressed the significant uncertainties in their result arising from the small sample size. A larger sample of 1800 galaxies drawn from the Sloan Digital Sky Survey (SDSS; Stoughton et al. 2002) was analysed by Tasca & White (2005) using the publicly available code GIM2D. They found that  $54 \pm 2\%$  of the local cosmic luminosity density at both  $r$  and  $i$  comes from disks and  $32 \pm 2\%$  from “pure bulge” systems and the remaining  $14 \pm 2\%$  from bulges in galaxies with detectable disks.

In this paper, we perform spheroid/disk decompositions from  $r$ -band images of a much larger sample of galaxies in the SDSS. The structure of the paper is as follows. In §2, we describe our dataset and how it is processed. In §3, we present results from the spheroid/disk decomposition. In §4, we derive the luminosity functions of disks and spheroids, and present stellar mass functions and also the supermassive black hole mass function. Finally, in §5, we give our conclusions. Appendices describe extensive tests of the reliability of both GALACTICA and GIM2D (Appendix A), technicalities of the fitting process (Appendix B) and comparisons of our S/T ratio morphologies with more traditional morphological measures (Appendix C). A cosmological model with  $\Omega_0 = 0.3$ ,  $\Lambda_0 = 0.7$  is adopted throughout and the Hubble constant is defined to be  $H_0 = 100 h \text{ km s}^{-1} \text{ Mpc}^{-1}$ .

## 2 DATA: SLOAN DIGITAL SKY SURVEY

### 2.1 Basic properties

The Sloan Digital Sky Survey (SDSS) is the largest imaging and spectroscopic survey to date. The SDSS Early Data

Release (EDR), made publicly available in 2001, consists of a 462 square degree area imaged in five pass-bands ( $u, g, r, i$  and  $z$ ) and also covered spectroscopically. The SDSS EDR galaxy catalogue is spectroscopically complete down to  $r = 17.7$  and contains measurements of various galaxy parameters (Stoughton et al. 2002). The imaging data were taken with a dedicated 2.5m telescope in the drift-scan (time-delay) integration mode with an effective exposure time of 54s. The data used in this study are the  $r$ -band imaging frames with corrections for bias, flat field, cosmic ray and pixel defects (Lupton et al. 2001). Each imaging frame is a  $2048 \times 1489$  pixel array with a pixel size of  $0.394''$ .

### 2.2 SDSS apparent magnitude limit

Benson et al. (2002) measured spheroid-to-total (S/T) light ratios for the field galaxy sample of Gardner et al. (1996) using  $I$ -band imaging. The data were originally obtained to determine the  $K$ -band luminosity function and Benson et al. (2002) showed that they could be used reliably to estimate S/T ratios for galaxies brighter than  $I_{\text{Gar}} = 16.0$  with an rms accuracy of  $\sigma_{\text{rms}} \sim 0.1$ . Unfortunately, the area covered by this sample is rather small ( $4.4 \text{ deg}^2$ ).

The SDSS imaging data were obtained with a larger telescope but using shorter exposure times than those of Gardner et al. (1996). From Monte-Carlo simulations, Benson et al. (2002) established the signal-to-noise required to obtain reliable measurements of S/T using the GALACTICA decomposition code. Assuming that, when applied to the SDSS data, the code will be reliable to the same overall signal-to-noise level, we find that the limiting magnitude required for our SDSS sample is  $I_{\text{SDSS}} - I_{\text{Gar}} = 0.4$ . Using the mean galaxy colours of Fukugita et al. (1995), the transformation between the  $I_{\text{SDSS}}$  and  $r$  bands is  $r - I_{\text{SDSS}} = 0.9$ , making the total difference equal to  $r - I_{\text{Gar}} = 1.3$  magnitudes. We therefore select EDR galaxies with  $r \leq 17.3$ .

### 2.3 SDSS data selection and galaxy catalogue

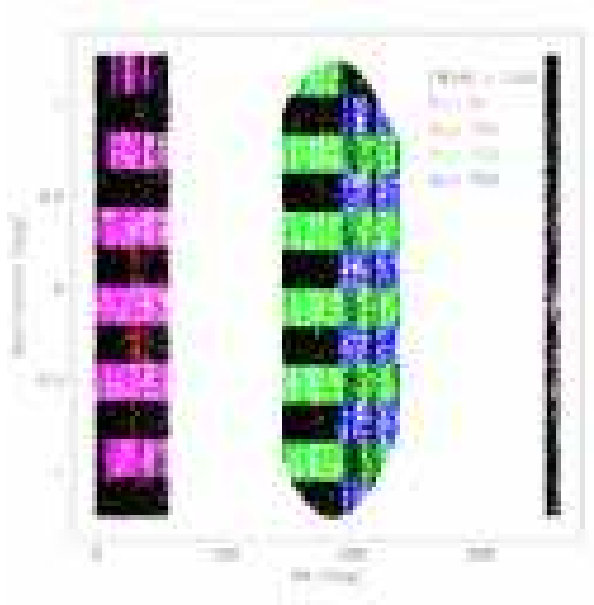
Galaxies with  $r \leq 17.3$  in the SDSS EDR equatorial strip are plotted in Fig. 1, colour-coded according to the SDSS run number (94, 125, 752 and 756). The black points represent imaging taken in ‘poor’ seeing conditions ( $\text{PSF}_{\text{FWHM}} > 1.55''$ , where  $\text{PSF}_{\text{FWHM}}$  is the full-width at half-maximum of the point spread function). This cut on of the seeing is used to impose a second galaxy selection criterion since reliable spheroid-to-disk decompositions require that the seeing be less than a typical galaxy half-light radius (Beijersbergen et al. 1999).

The final galaxy selection criterion is redshift. To avoid contamination of the measured redshift by the local galaxy infall velocity, a low redshift cut,  $z = 0.02$ , is imposed. Since the total SDSS sample begins to tail off at large distance, a high redshift cut,  $z = 0.3$ , was also imposed.

The selection leads to a total of 8839 SDSS EDR galaxies. The relevant SDSS pipeline parameters extracted from the catalogue of (Lupton et al. 2001) are summarized in Table 1.

**Table 1.** A sample of the main catalogue parameters. Column 1: galaxy and run ID—necessary to identify the CCD frames that contain galaxy/galaxies of interest. Columns 2 & 3: right ascension and declination (degrees)—used to identify the catalogued galaxies within a given CCD frame. Column 4: Reddening corrected  $r$ -band Petrosian magnitude. Column 5:  $r$ -band Petrosian radius,  $R_P$ , in arcseconds, as defined by Stoughton et al. (2002)—these radii are used to define the postage-stamp size for a given galaxy (see Appendix A2.1). Column 6:  $r$ -band seeing (the mean value of stellar FWHM in arcseconds)—the GALACTICA code uses this value as the initial estimate of the seeing parameter. Column 7: redshift—used for computing absolute magnitudes

ID	RA (deg)	Dec (deg)	$r_P - r_{\text{redd}}$	$R_P$ (")	FWHM (")	z
100	2.6311	-0.0429	17.24	8.9	1.5	0.039
101	2.6639	-0.0529	14.36	14.9	1.5	0.039
102	2.3471	-0.0561	16.99	4.3	1.4	0.076
103	2.3398	-0.0704	16.68	11.9	1.4	0.058
104	2.2699	-0.0005	16.90	8.0	1.4	0.085
105	2.2338	-0.1939	16.31	6.8	1.4	0.138
106	2.3138	-0.1303	16.87	4.8	1.4	0.116
107	2.1515	-0.0513	16.48	9.3	1.4	0.076
108	2.0482	-0.0493	17.09	4.3	1.4	0.101
109	2.0822	-0.0007	17.20	3.9	1.4	0.079
110	2.1158	-0.0005	16.99	3.3	1.4	0.158



**Figure 1.** SDSS galaxies that meet our selection criteria. Black points correspond to galaxies imaged when the seeing was greater than  $1.55''$ .

## 2.4 Sample solid angle

To calculate the solid angle covered by our sample, the galaxy coordinates were accumulated in  $0.2^\circ$  bins. All the areas which contain at least one galaxy residing in a particular bin were summed to give the total solid angle. The bin size was chosen such that the derived solid angle was insensitive to small changes in the bin size. As an additional check, this same bin size was used to reproduce the SA of the entire SDSS EDR. For our chosen bin size, the solid angle subtended by our sample is  $165.5$  square degrees.

## 2.5 Object detection and astrometry

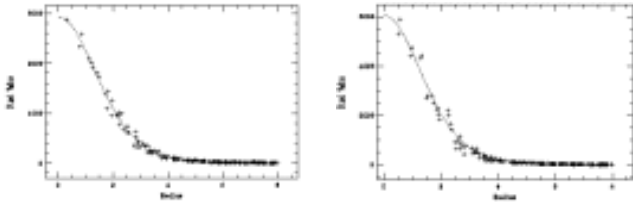
Object detection was performed using SExtractor v2.2.2 (Bertin & Arnouts 1996). The SExtractor world coordinates of the object centroid positions ( $x, y$ ) were used to identify the catalogued galaxies within the SDSS frames. The GALACTICA code (Appendix A2.3) was run on the extracted postage stamps whose size was set equal to  $(2 \times R_p) \times (2 \times R_p)$ , where  $R_p$  is the Petrosian radius (Lupton et al. 2001). This is large enough to contain many background pixels but sufficiently small to ensure a reasonable convergence time for the fitting procedure. Prior to decomposition, the SExtractor estimate of the local sky background was subtracted from every postage-stamp to ensure that the background level was close to zero (see Appendix A).

## 2.6 SDSS point spread function (PSF)

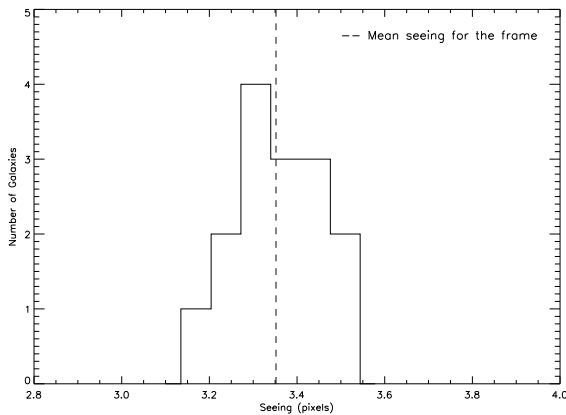
Before beginning the decomposition procedure, it is necessary to ensure that the PSF analytic model assumed by the GALACTICA code (Appendix A2.3) is a realistic representation of the SDSS PSF. To demonstrate that the SDSS stars are well represented by the analytic Moffat profile assumed by the GALACTICA code, the IRAF task IMEXAMINE was used to fit radial Moffat profiles to a sample of stars imaged on different SDSS frames and at different positions within every frame. Fig. 2 shows radial fits to stellar light profiles obtained using  $\beta = 4.5$  and demonstrates that a Moffat star with this value of  $\beta$  is a good analytic representation of the SDSS PSF.

### 2.6.1 PSF variation

The GALACTICA code assumes a starting value for the PSF equal to the measured value of the seeing in the SDSS and allows the value to fluctuate by  $\pm 5\%$  (Appendix A2.3). The  $\pm 5\%$  variation is set from the observed variation of the seeing across a typical SDSS frame, as shown in Fig. 3, which demonstrates that for stars imaged at various positions in a given SDSS frame, the FWHM does not change by more than  $\pm 5\%$ .



**Figure 2.** The radial Moffat profile ( $\beta = 4.5$ ) fits (solid line) to the SDSS stellar light profiles (points) for stars found at various positions within several SDSS frames. The radius is in pixels and the Pixel Value are counts. Results are shown for two different frames.



**Figure 3.** The variation of the PSF across an SDSS frame. The dashed line represents the mean value of the seeing for the frame. The seeing appears not to vary by more than  $\pm 5\%$  from the mean value. A similar inspection of other frames shows this to be true in general.

The small allowed change in the seeing ensures that the GALACTICA code can find the representative value of the seeing at each galaxy position. However, it is important to test how consistently the GALACTICA code recovers the ‘correct’ representative PSF for a given galaxy and quantify the effect this has on the recovered S/T ratios. The observed galaxy properties are expected to vary little between the  $r$  and  $i$  bands but the PSF signatures for these observations will be somewhat different. Fig. 4 shows a good correlation between the S/T ratios obtained for the same set of galaxies imaged in the two bands. The code finds consistent S/T ratios across a range of apparent magnitudes independently of the seeing. Error bars are obtained from 30 Monte Carlo realizations of each of the model fits, assuming the noise appropriate to each image.

### 3 GALACTICA DECOMPOSITIONS AND GALAXY MORPHOLOGIES

#### 3.1 GALACTICA decomposition outputs

Fig. 5 demonstrates a typical fit to a galaxy light profile. The figure shows the postage-stamps of a real galaxy, a noise-free model generated from the best-fit parameters along with the individual model disk and spheroid components. In this

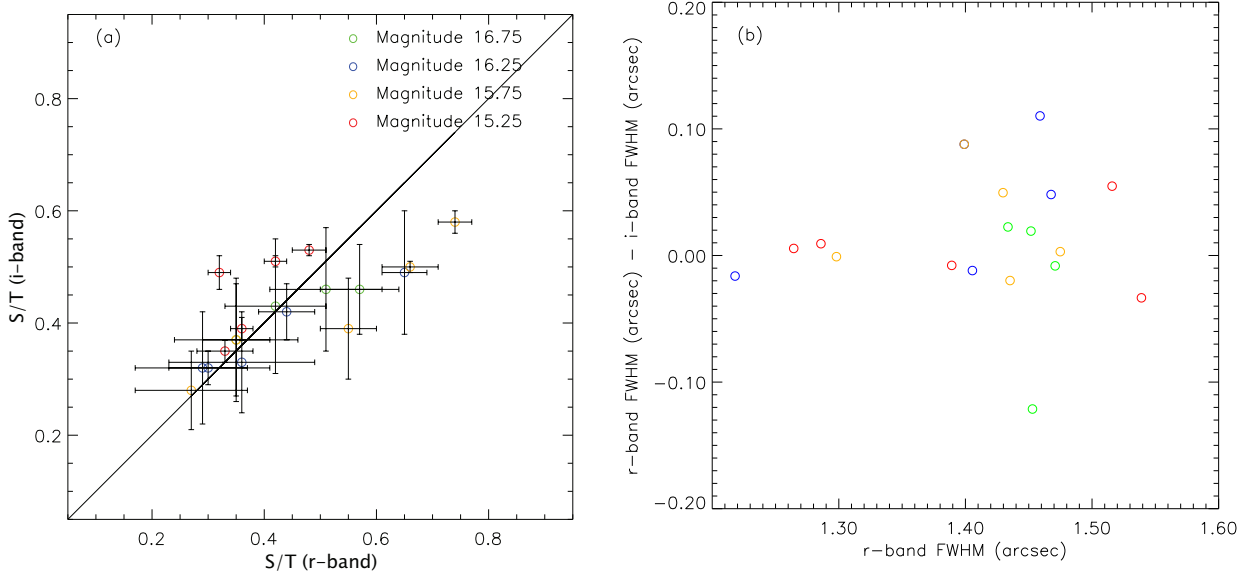
$$S/T = 0.30 \pm 0.07$$

**Figure 5.** Top: real (left) and model (right) images. Bottom: disk (left) and spheroid (right) component fits. The cross-hatched regions represent potential contamination from overlapping objects (or regions where data were unavailable after the image was re-centred by GALACTICA) and are excluded from the fitting. The contours indicate the pixel values in ADU/s.

$$S/T = 0.30 \pm 0.07$$

**Figure 6.** Real (left) and residual (right) images. The inset shows the distribution of S/T ratios from 30 Monte Carlo realisations with the vertical dashed line indicating the best-fit S/T. The value of  $\chi^2_\nu$  is acceptably small and the residual image also shows a good fit to the data.

study, a galaxy is deemed to be sufficiently well represented by the model if the  $\chi^2$  per degree of freedom satisfies  $\chi^2_\nu < 2.0$ , and if there are no obvious structures left in the residual image. An example of a well fit galaxy is shown in Fig. 6. The cross-hatched areas represent potential contamination from overlapping objects as determined by the GALACTICA masking procedure (see Appendix A2.3) and are excluded from the fitting. The inset in Fig. 6 shows a histogram of  $dP/d(S/T)$  —the distribution of the spheroid-to-total ratio from 30 Monte Carlo realisations, with the vertical dashed line indicating the best-fit S/T value for this galaxy.



**Figure 4.** (a) Correlation between S/T ratios obtained for the same galaxies in the  $r$  and  $i$  bands. The good correlation demonstrates that the S/T ratios are accurately determined for different representative PSFs and across the apparent magnitude range. The recovered S/T ratios show no obvious dependence on galaxy apparent magnitude, indicating that the decompositions are not affected by the variation in the signal-to-noise ratio. (b) The difference in the output GALACTICA seeing for the same set of galaxies observed in the  $r$  and  $i$  bands. The lack of a trend demonstrates that the GALACTICA code recovers the representative PSF for each galaxy well, without biasing the recovered S/T ratios.

### 3.2 Correlations of S/T with other fit parameters

Understanding the properties of this large statistical sample of galaxies is important since it may reveal features which otherwise would not be discovered in smaller samples such as those discussed in Appendix C. Equally, any unexpected correlations between parameters could help discover and reduce possible biases introduced by the fitting routine.

Histograms of various properties of our SDSS galaxies inferred from the GALACTICA decompositions are shown in Fig. 7. These plots reveal the following:

- (i) a large number of highly elliptical spheroids;
- (ii) an excess in the number of galaxies with spheroid position angle,  $\theta_s$ , equal to  $0^\circ$  and  $180^\circ$ ; and,
- (iii) a non-uniform distribution of the cosine of the disks' inclination,  $\cos(i)$ .

In the remainder of this section, we explore the possible origins of these unexpected distributions and their influence on the recovered values of the S/T ratio.

#### 3.2.1 S/T vs. Ellipticity

Around 15% of galaxies appear to have a highly elliptical spheroid component whose ellipticity has reached the imposed upper limit<sup>2</sup> of  $e = 0.83$ . A large number of frames have been inspected by eye and show that these galaxies generally exhibit bar-like structures in the direction of the

detected highly elongated spheroidal component. In these cases, the existence of this extra component, which is not part of the fitted model, drives the code to fit small and highly elliptical spheroids (Fig. 8). These galaxies are generally disk-dominated and thus we expect this shortcoming of the model to introduce only a small bias on the overall S/T ratio. Any bias that is introduced would increase this ratio, resulting in a slight overestimation of the spheroid luminosity density in §4.

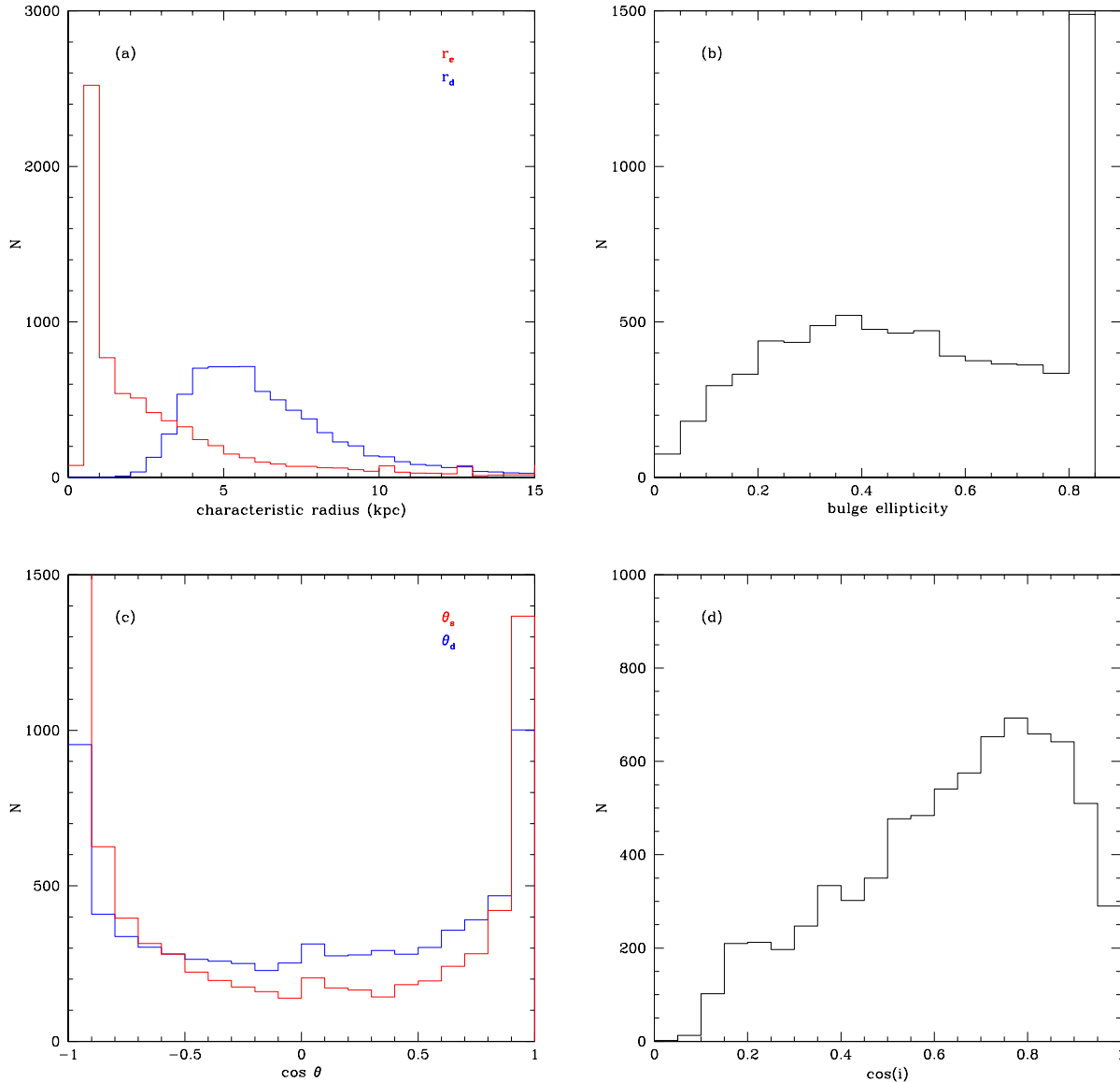
#### 3.2.2 S/T vs. spheroid position angle

Many galaxies appear to have  $\theta_s \sim 0^\circ$ . This could be due to either:

- (i) some feature intrinsic to the code such as the initial estimate of  $\theta_s$ ; or,
- (ii) a feature intrinsic to the data.

For our purposes, the crucial issue is whether the bias in  $\theta_s$  affects the derived S/T. To test (i), a galaxy with the recovered  $\theta_s \sim 0^\circ$  was re-fitted three times. Each time, the initial value of  $\theta_s$ , estimated directly from the image, was ignored and an initial (and very different)  $\theta_s$  was specified. The code was found to converge very quickly to a minimum at  $\theta_s \sim 0^\circ$  with the same S/T ratio irrespective of the very different initial  $\theta_s$ , demonstrating that the initial value of  $\theta_s$  is not important. Point (ii) is a plausible explanation since the data were taken in drift-scan mode along the easterly direction which corresponds to  $\theta = 0^\circ$ . To quantify how dependent is the derived S/T upon the spheroid orientation, we randomly selected a sample of 100 SDSS galaxies, rotated the images by  $+90^\circ$  and re-fitted. The results, displayed in

<sup>2</sup> The upper limit for the ellipticity corresponds approximately to that of the most elliptical observed galaxies (Lambas, Maddox & Loveday 1992).

$\chi^2 < 2$ ; 7493/8839 galaxies

**Figure 7.** Histograms of parameters recovered by GALACTICA from our sample of SDSS galaxies. a) Characteristic sizes of the spheroidal and disk components. b) Ellipticity of the spheroidal components. c) Cosine of the position angles of disks and spheroids. d) Cosine of the inclination angle of the disks. Where appropriate, spheroids are represented in red and disks in blue.

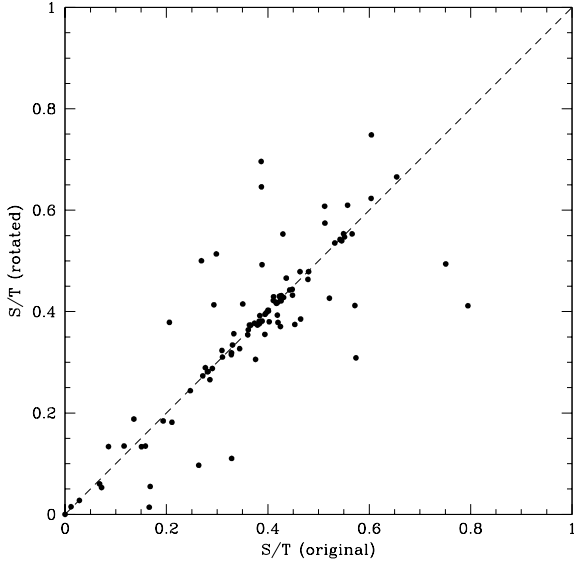
Fig. 9, indicated that the recovered S/T ratios ( $\sigma_{\text{rms}} = 0.09$ ) are not strongly affected by the orientation of the spheroid position angle on the sky.

### 3.2.3 S/T vs. disk inclination

A large number of objects in the sky which are randomly inclined to the line-of-sight should have a uniform distribution of  $\cos(i)$ . Fig. 7 clearly shows that this is not the case for the inclination angles of the disk components obtained by decomposing our sample of SDSS galaxies.

To test whether the apparently incorrect recovery of the disk inclination is an artifact of the fitting procedure, a sample of 200 mock galaxies was generated using the GALAC-

TICA code (see Appendix A). The S/T ratios were chosen at random in the interval  $[0, 1]$ . The remaining parameters, including the value of  $\cos(i)$ , were also chosen at random. Fig. 10 demonstrates that the GALACTICA code generally recovers the  $\cos(i)$  distribution for 200 model galaxies quite well. However, a noticeable feature is a slight excess around  $i > 75 - 80^\circ$ . This excess reflects that fits avoid the  $90^\circ$  limit since this would correspond to fitting an infinitely thin edge-on disk and, because of seeing, the disks are never infinitely thin edge-on. (The feature remains even if the allowed inclination range is increased from  $[0^\circ, 90^\circ]$  to  $[-180^\circ, 180^\circ]$ ). The S/T ratios are not affected by this problem, i.e. model galaxies with input value  $i \sim 90^\circ$  but recovered value  $i \sim 85^\circ$  show no bias in the recovered S/T ratio.



**Figure 9.** *Left panel:* Comparison of estimates of  $S/T$  for a sample of 100 SDSS galaxies estimated from the actual image and from the image rotated anticlockwise by  $90^\circ$ . *Right panel:* The difference in the recovered position angle before and after rotation. The recovered  $S/T$  ratios are not strongly affected by the orientation of the spheroid’s position angle.

The questions that remain to be answered are:

(i) what is the reason for the  $\cos(i)$  bias, and (ii) how well is the  $S/T$  ratio recovered if the disk inclination angle is recovered incorrectly?

To answer the first question, we have examined in detail the  $\cos(i)$  distribution for sub-samples of our galaxies selected in different ways. We find that the strongest variation in the  $\cos(i)$  distribution is seen when we select galaxies on the basis of the angular size of their disk component relative to the seeing. The upper panel of Fig. 11 shows the distribution of  $\cos(i)$  for various sub-samples selected by disk angular size,  $a_{\text{disk}}$ . It is clear from Fig. 11 that the galaxies with the most biased  $\cos(i)$  distribution are those with  $a_{\text{disk}} < 2\sigma_{\text{PSF}}$ , where  $\sigma_{\text{PSF}}$  is defined in Appendix A2.3, that is, those galaxies which have a very small angular size compared to the seeing. Conversely, for the sample with  $a_{\text{disk}} > 6\sigma_{\text{PSF}}$  we find a much more uniform  $\cos(i)$  distribution.

Such an effect is not surprising: if the disk is poorly resolved (i.e. small relative to the PSF), there will be little constraint on its inclination. In particular, a highly inclined disk will be made to appear more circular (and hence less inclined) by the effects of seeing. It is important to note that GALACTICA (and GIM2D also) actually use  $i$  as a fitting parameter, rather than  $\cos(i)$ . The lower panel of Fig. 11 shows the distribution of  $i$  for our whole galaxy sample. As shown by the line in this panel, the distribution of  $i$  has an approximately Gaussian shape, a clear signal that  $i$  is essentially unconstrained for the majority of our galaxies.

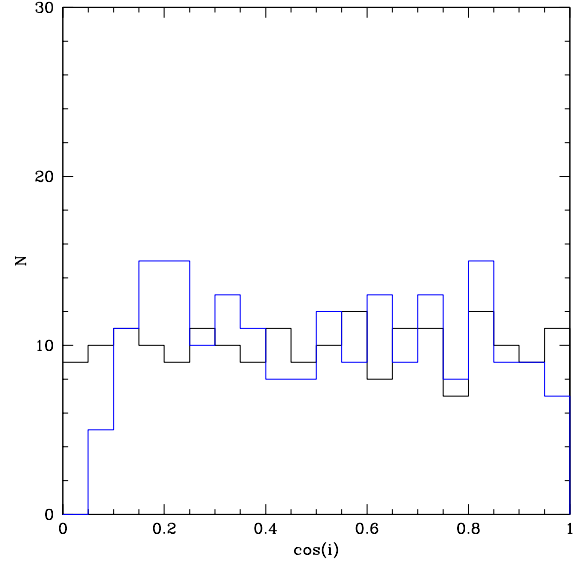
There remains, however, a lack of galaxies with  $\cos(i) \lesssim 0.2$  even in our well-resolved sub-sample. As noted in Appendix A3.2, in tests with mock galaxies the GALACTICA code tends never to fit an inclination above  $i \approx 85^\circ$  ( $\cos i \approx 0.09$ ) even for edge-on disks (a similar effect is found in GIM2D). This occurs because, even for an edge-on disk, the

finite pixel size and PSF will artificially “thicken” the disk image, allowing a somewhat less inclined model to be fit to the image. However, this alone cannot explain the lack of galaxies with  $\cos(i) \lesssim 0.2$ . The solid line in Fig. 11 shows the distribution of  $\cos(i)$  expected if infinitely thin disk models are fitted to mock galaxies which possess disks of finite thickness. Specifically, in this case we have assumed a distribution of disk thicknesses,  $h = z_0/r_d$ , where  $z_0$  is the vertical scale height of the disk, which is log-normal, with mean  $h = 0.3$  and dispersion  $\sigma_{\ln h} = 0.5$ , motivated by the calculations of Benson et al. (2004). This distribution of thickened disks reproduces the  $\cos(i)$  distribution for our well-resolved sub-sample quite well.

Such biases in the distribution of  $\cos(i)$  have been seen in other studies employing 2D galaxy decomposition techniques. Tasca & White (2006) used GIM2D to fit the 2D images of galaxies in the SDSS. They found a biased distribution of  $\cos(i)$ , with intrinsically brighter galaxies showing the most biased distribution. Fig. 12 reproduces Fig. 10 of Tasca & White (2006), with results from this work overlaid. Our results, using the same dataset but a different galaxy decomposition code, are in excellent agreement with those of Tasca & White (2006). This variation of the  $\cos(i)$  distribution with intrinsic luminosity is to be expected: more luminous galaxies can be seen at greater distances and so may be expected to have, on average, smaller angular sizes. As a result, more luminous galaxies will typically be less well resolved than fainter ones, and so their inclinations are less well constrained.

Allen et al. (2006) performed 2D galaxy decompositions, also using GIM2D, on galaxies in the Millennium Galaxy Catalogue and found that disk-dominated galaxies ( $S/T < 0.8$ ) had a more uniform (although still biased) distribution of  $\cos(i)$ . Fig. 13 reproduces their results, with comparable results from our own work overlaid. In this





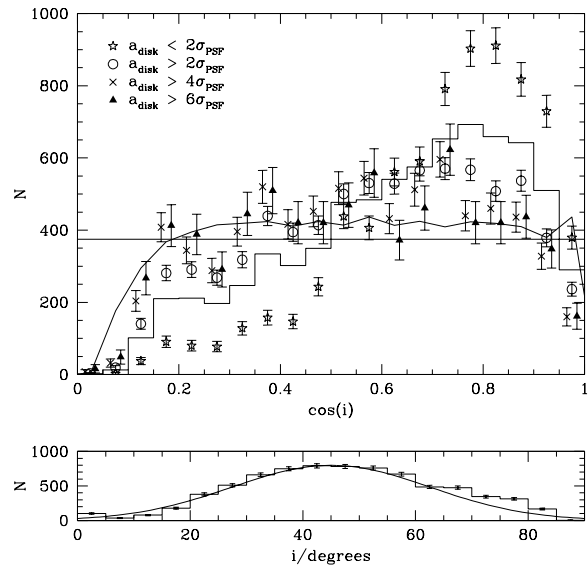
**Figure 10.** The input (black) and the recovered (blue)  $\cos(i)$  distribution for 200 model galaxies created and decomposed using the GALACTICA code. The figure demonstrates that the non-uniformity in the  $\cos(i)$  is not caused by the fitting code. An apparent excess of galaxies with  $i \sim 75 - 80^\circ$  can be seen.

$$S/T = 0.37$$

**Figure 8.** An example of a galaxy with a highly elliptical spheroid. This galaxy demonstrates how the central bar-like structure in the galaxy results in the detection of a highly elliptical spheroid along the same direction. The top row shows the real (left) and model (right) images. The middle row shows the disk (left) and spheroidal (right) components. The bottom row shows the real (left) and residual (right) images.

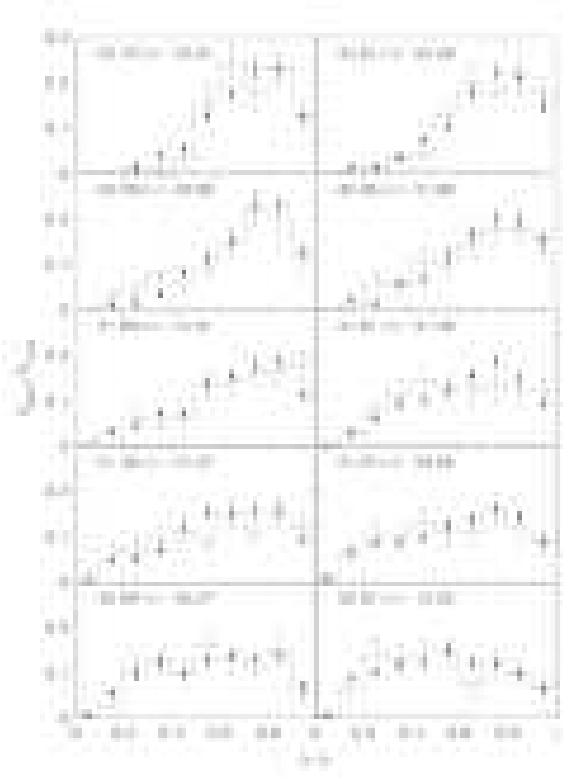
case, we find the opposite trend: our  $\cos(i)$  distribution is more uniform for the  $S/T > 0.8$  sample, although the errors are large. We find that galaxies must have angular sizes of several times the seeing half-width at half-maximum in order for the inclination to be well constrained. From Fig. 1 of Allen et al. (2006), we would therefore conclude that a large fraction of their galaxies should have poorly constrained inclinations.

In conclusion, we find no evidence for a true *bias* in  $\cos i$ , merely an inability to constrain it for galaxies which are poorly resolved. Likewise, we find that the  $S/T$  ratios of galaxies are poorly constrained when the galaxy in question is poorly resolved, *unless* the galaxy happens to be seen almost edge-on which permits a better discrimination between disk and spheroid even in poorly resolved galaxies. Over 50% of our galaxies are either sufficiently well resolved, or seen sufficiently edge-on to have well constrained  $S/T$  ratios. If we excise galaxies with poorly constrained  $S/T$  ratio from our sample we find that the resulting luminosity functions of disks and spheroids are not strongly affected—we find a 50% increase in abundance of the faintest spheroids and a slight reduction in the characteristic luminosity  $L_*$ . These effects



**Figure 11.** *Upper panel:* The distribution of  $\cos i$  split by the angular size of galaxy disks. The histogram shows the distribution for our complete galaxy sample, while the horizontal line indicates that expected for a randomly oriented set of galaxies. Symbols indicate the distribution for sub-samples selected by disk angular size ( $a_{\text{disk}}$ ; see the legend for selections). The solid line indicates the expected distribution of  $\cos i$  if a perfectly thin galaxy disk is fit to galaxy disks which are in fact thick, with a log-normal distribution of disk thicknesses with mean  $h = 0.3$  and dispersion  $\sigma_{\ln h} = 0.5$ . *Lower panel:* The distribution of inclination angle,  $i$ , for our galaxy sample is shown by the histogram. As indicated by the curve, this approximates a Gaussian distribution centred on  $45^\circ$ .





**Figure 12.** The distribution of axial ratio  $b/a$  (equivalent to  $\cos i$ ) as a function of galaxy luminosity. Histograms show results from Tasca & White (2006) (this figure is a reproduction of their Fig. 10), while points show results from this work. Both datasets show that the bias in  $\cos i$  occurs primarily for the most luminous galaxies.

are consistent with the random scatter in luminosities introduced into galaxies for which  $S/T$  is poorly constrained.

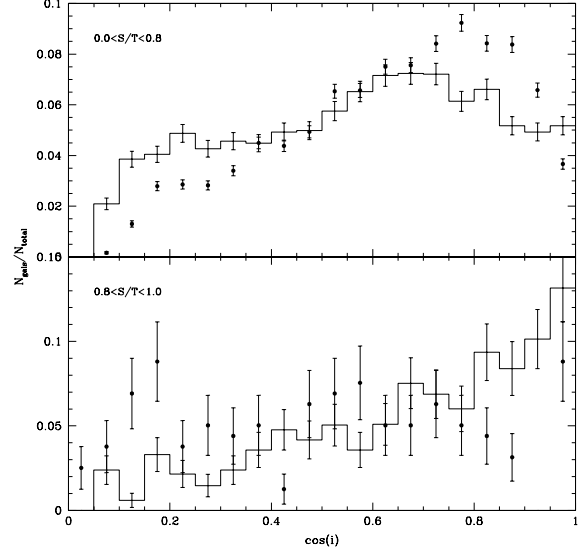
## 4 LUMINOSITY AND MASS FUNCTIONS

### 4.1 Introduction

The spatial abundance of galaxies is expressed by means of the luminosity function (LF), defined as:

$$dn(M) = \phi(M)dM, \quad (1)$$

where  $dn$  is the number density of galaxies with absolute magnitude in the range  $M, M + dM$ . The simplest way to calculate the luminosity function is using the  $1/V_{\max}$  method in which the number of galaxies in each individual absolute magnitude bin is divided by the volume of space that has been surveyed at that magnitude. Galaxies in any given absolute magnitude range are assumed to be uniformly distributed in the surveyed volume which is not the case if any local overdensities are present. Maximum likelihood techniques circumvent this problem and provide more accurate estimates of the luminosity function. Here, we will employ the  $1/V_{\max}$  as well as the Stepwise Maximum Likelihood (SWML) non-parametric estimator (Efstathiou, Ellis & Peterson 1988) which characterizes the LF as a series of steps. We will also employ the STY parametric estimator Sandage,



**Figure 13.** The distribution of  $\cos i$  split by  $S/T$ . Histograms show results from Allen et al. (2006) (this figure is a reproduction of their Fig. 9), while points show results from this work. While both datasets show biased distributions of  $\cos i$ , the trends with  $S/T$  appear to differ, with the Allen et al. (2006) dataset showing a more uniform distribution of  $\cos i$  for galaxies with low  $S/T$ .

Tammann & Yahil (1979), assuming a Schechter functional form: (Schechter 1976):

$$\phi(M) = 0.4 \ln 10 \phi_* 10^{-0.4(M-M_*)(\alpha+1)} \exp[-10^{-0.4(M-M_*)}], \quad (2)$$

where  $M_*$  is a characteristic magnitude,  $\alpha$  is the faint-end slope and  $\phi_*$  is the normalization. Integrating over the Schechter function provides an estimate of the luminosity density. This can also be obtained by summing up all the individual SWML contributions.

Computing the spheroid and disk LFs is more complicated since there is an additional constraint to be considered (Benson et al. 2002) namely the detectability of a spheroid/disk depends both on the component apparent magnitude and the corresponding  $S/T$ . This needs to be accounted for when constructing the luminosity function. A detailed discussion of the application of these methods can be found in Benson et al. (2002). Benson et al. (2002) used a functional form for STY parametric fits to the spheroid and disk LFs which had a Schechter $\times$ exponential form. We find that Benson et al.'s functional form does not provide a good description of our larger sample of galaxies. We have been unable to find a suitable functional form which does provide a good description and so have not performed STY fits to the spheroid and disk luminosity function data.

### 4.2 SDSS absolute magnitudes and K+E corrections

In order to estimate the luminosity function, we require galaxy absolute magnitudes. A galaxy at redshift  $z$ , with apparent magnitude  $m$ , has an absolute magnitude  $M$  given by:

$$m - M = 25 + 5 \log_{10}(D_L) + KE(z) \quad (3)$$

where  $D_L$  is the luminosity distance in megaparsecs and  $KE(z)$  is the K+E correction.

K+E corrections for our catalogued galaxies were obtained using a code kindly provided by Carlton Baugh. It employs the revised isochrone stellar population synthesis models of Bruzual & Charlot (1993) to determine present-day galaxy luminosities. The model assumes a stellar initial mass function (IMF) and a star-formation rate  $\psi(t) \propto \exp(-t/\tau)$ , with timescale,  $\tau$ . A grid of models was generated by varying the metallicity and  $\tau$ . We assume a Salpeter (1955) IMF and apply a simple dust extinction law. At every point on the grid, a table of absolute magnitudes, galaxy colours, K+E corrections and galaxy mass-to-light ratio is generated. The model that best matches the observed  $g-r$  and  $r-i$  colours of each galaxy is then used to infer its present-day ( $z=0$ )  $r$ -band absolute luminosity, K+E correction and stellar mass-to-light ratio. The mass-to-light ratio is used to convert luminosities to stellar masses in order to estimate stellar mass functions (see §4.5). Note that the K+E corrections are based on the total (i.e. spheroid plus disk) colour of a galaxy.

### 4.3 Luminosity function estimates

Using the methods developed by Benson et al. (2002), we estimate the luminosity functions of spheroids and disks, as well as the total galaxy luminosity function for our sample of SDSS EDR galaxies. Our results are displayed in Fig. 14. For galaxies whose images are well-fit by our model, we find that the STY method accurately recovers the parameters of the total luminosity function; furthermore, the STY fit traces the corresponding SWML points very well. The value of  $M_*$  obtained from the STY fit to the total luminosity function agrees very well with that of Nakamura et al. (2003) (SDSS  $r$ -band,  $z=0$ ), although our estimate of  $\alpha$  is different from theirs at the  $2.5\sigma$  level. While we have not been able to find a parametric form which fits to spheroid and disk 2D luminosity functions ( $\Phi(M, S/T)$ ) we have determined the parameters of Schechter functions which fit the SWML data points reasonably well. These should not be considered good fits in a statistical sense, merely useful fitting functions. The parameters of the best fit Schechter functions are given in Table 2.

We calculate luminosity densities of disks and spheroids by integrating over the SWML points<sup>3</sup>. We find the luminosity densities for spheroids and disks to be:  $\rho_L = 0.62 \pm 0.08 h L_\odot \text{ Mpc}^{-3}$  and  $\rho_L = 1.84 \pm 0.27 h L_\odot \text{ Mpc}^{-3}$  respectively. This is in contrast with the findings of the previous study of Benson et al. (2002) who found the spheroid and disk luminosity densities to be very nearly equal. Of course, Benson et al. (2002) used a very small sample of galaxies to compute luminosity densities, finding a ratio of disk to spheroid luminosity density of  $1.2 \pm 0.9$ . Our current

sample gives a ratio of  $2.97 \pm 0.58$ , which is consistent with that of Benson et al. (2002) at the  $2\sigma$  level.

#### 4.3.1 Effects of S/T bias on luminosity functions

In Appendix A3.2 we find that the value of S/T recovered by GALACTICA (and also GIM2D) for mock images are biased. In this section, we address how this affects our estimates of luminosity functions. The median bias in S/T produced by GALACTICA can be approximated by a linear dependence on the true S/T (see Fig. A6). We use this linear relation to apply a correction to the value of S/T recovered for each SDSS galaxy in order to obtain an estimate of the unbiased value. We then construct new disk and spheroid luminosity functions using these corrected values.

We find that there are only small changes in the measured luminosity functions, the most significant being a small enhancement in the abundance of bright spheroids. The luminosity density ratio quoted above varies by less than  $0.5\sigma$  after correcting for this bias.

### 4.4 Comparison with theoretical predictions

In Fig. 15 we compare our estimate of the disk and spheroid luminosity functions with predictions from the Baugh et al. (2005) and Bower et al. (2006) implementations of the GALFORM semi-analytic model of galaxy formation. These two models differ in a number of important respects. For example, in Bower et al. (2006) feedback from the emission of active galactic nuclei plays a role in quenching cooling flows in clusters; in the Baugh et al. (2005) model, a top-heavy IMF is assumed for stars that form in starbursts. The two models, however, assume similar mechanisms for the formation of disks and spheroids: disks form when spinning gas cools in a halo while spheroids form either by major mergers or by instabilities in the disks. Although both models generally provide a reasonable description of many galaxy properties, they have different strengths and weaknesses. Neither of them has been previously applied to the study of galaxy morphology.

Fig. 15 shows that both GALFORM models reproduce the main trends seen in the SDSS luminosity functions. At faint magnitudes, the luminosity function is dominated by disks while at bright magnitudes disks and spheroids make comparable contributions. The Bower et al. (2006) model in particular provides a good match to the SDSS luminosity functions.

### 4.5 Stellar mass functions

As noted in §4.2, our procedure for determining K+E corrections also provides an estimate of the stellar mass of each galaxy. Using these stellar masses we have constructed total, spheroid and disk stellar mass functions using the SWML and STY (for total mass only) methods. For the spheroid and disk mass functions we also show the Schechter function which best fits the SWML data points. The results are shown in Fig. 16, with parameters of Schechter function fits shown in Table 3. It should be noted that we implicitly assume that the mass-to-light,  $Y$ , ratio determined for each by our K+E correction procedure is the same for both the disk

<sup>3</sup> No correction is included for galaxies fainter than the lower limit shown in the figures. Using the best-fit Schechter functions listed in Table 2 we estimate that including fainter spheroids/disks would lead to corrections of 4%/40% respectively. We regard these corrections as speculative since the Schechter function does not provide a good fit to the spheroid and disk luminosity functions.

**Table 2.** Best fitting Schechter function parameters for luminosity functions of total, spheroid and disk components. For the total luminosity function the best fit parameters are determined using the STY method. For the disk and spheroid luminosity function we instead fit a Schechter function to the non-parametric luminosity function determined using the SWML method—these should be considered useful fitting functions only, not good fits in any statistical sense. For the spheroid and disk luminosity function fits, the maximum deviation from the SWML data points is given in the final column.

Component	$M_* - 5 \log h$	$\alpha$	$\phi_0/h^3 \text{Mpc}^{-3}$	Max. Dev.
Total	-20.87	-1.57	0.0092	N/A
Spheroid	-20.91	-1.40	0.0023	20%
Disk	-20.61	-1.74	0.0065	23%

and spheroid components. In reality, the recovered value of  $\Upsilon$  reflects some weighted average of the  $\Upsilon$  of each component. To improve upon this situation would require a more advanced procedure in which the  $\Upsilon$  (and K+E correction) of disk and spheroid components were estimated separately using measurements of the disk and spheroid colours. This would require performing spheroid-disk decompositions in multiple wavebands.

A quantity of interest is the ratio of stellar mass in disks to that in spheroids, averaged over the entire galaxy population. Integrating the stellar mass densities displayed in Fig. 16 we obtain the average density of stars in disks and spheroids in units of the critical density. We find  $\Omega_{\text{stars,disks}} = (0.656 \pm 0.005)h^{-1}10^{-3}$  and  $\Omega_{\text{stars,spheroids}} = (0.383 \pm 0.005)h^{-1}10^{-3}$ . These results are in good agreement with those of Benson et al. (2002) who found  $\Omega_{\text{stars,disks}} = (0.51 \pm 0.08)h^{-1}10^{-3}$  and  $\Omega_{\text{stars,spheroids}} = (0.39 \pm 0.06)h^{-1}10^{-3}$ . We conclude that the fraction of stellar mass found in disks today is  $63 \pm 1\%$ . This is comparable to (although formally significantly different from) the fraction of  $55 \pm 2\%$  reported by Tasca & White (2005).

#### 4.6 Black hole mass function

In the last few years, it has been conclusively demonstrated that many galaxies possess central supermassive black holes and that their mass is strongly correlated with the properties of the galaxy’s spheroid such as its luminosity, stellar mass and velocity dispersion (Kormendy & Richstone 1995; Magorrian et al. 1998; Merritt & Ferrarese 2001; Marconi & Hunt 2003; Häring & Rix 2004). Although there is only direct evidence for these black holes in bright galaxies, it seems quite plausible that galaxies of all sizes have a central black hole (e.g. Malbon et al. 2006).

From the mass function of galactic spheroids determined in §4.5, assuming that all spheroids contain a supermassive black hole at their centre, we can estimate the mass function of supermassive black holes in the local Universe. We assume that the black hole mass is given by  $M_{\bullet}/M_{\odot} = 1.6 \times 10^8 [M_{\text{spheroid}}/10^{11} M_{\odot}]^{1.12}$  (Häring & Rix 2004) and ignore any scatter in this relation since an accurate determination of the black hole mass function would first require a deconvolution of the (uncharacterized) error distribution of spheroid masses.

The resulting black hole mass function is shown in Fig. 17. Parameters of the Schechter function which best fits the SWML data points are given in Table 3. Integrating this mass function gives a total black hole mass density in

the local universe of  $\rho_{\bullet} = (1.77 \pm 0.08) \times 10^5 h M_{\odot} \text{Mpc}^{-3}$ . This number has a much smaller error and is a factor of 2–3 smaller than other determinations (Yu & Tremaine 2002; Aller & Richstone 2002; Marconi et al. 2004; McLure & Dunlop 2004; Shankar et al. 2004) based on much smaller samples of galaxies. This discrepancy may be due in part to the fact that our estimated mass-to-light ratios for galaxies are an average of the mass-to-light ratios of disk and spheroid components. Since the spheroid is likely to have a larger mass-to-light ratio than the disk, this will cause us to underestimate the mass of the spheroid component and therefore of the black hole. Note also that we ignore any scatter in the black hole-spheroid mass relation. Assuming a scatter of 0.3 dex (Häring & Rix 2004) would increase our inferred black hole mass density to  $\rho_{\bullet} = (2.40 \pm 0.13) \times 10^5 h M_{\odot} \text{Mpc}^{-3}$ .

For comparison with our inferred black hole mass function, we show results from the recent model of Malbon et al. (2006) who incorporate a calculation of supermassive black hole growth into the GALFORM semi-analytic model of galaxy formation in a  $\Lambda$ CDM universe. The lines in Fig. 17 show their results for two different specific galaxy formation models. While the calculation based upon the parameters of Baugh et al. (2005) seems to match the abundance of the high-mass black holes quite well, neither model is able to reproduce the inferred low abundance of less massive black holes.

Before our results can be used to constrain such models strongly, it will be necessary to achieve a significantly better understanding of the uncertainties in the measured spheroid mass, and to perform the conversion from luminosity to stellar mass using a technique which accounts for the different stellar populations in the spheroid and disk.

## 5 SUMMARY AND CONCLUSIONS

We have used a sample of  $\sim 9000$  galaxies extracted from the Sloan Digital Sky Survey to estimate the spheroid and disk luminosity and stellar mass functions in the local universe using the GALACTICA code of Benson et al. (2002). The 2D model fits to the surface brightness have revealed a bias in the recovered disk inclination angle arising from the lack of strong constraints on this parameter for most galaxies. Extensive tests have shown that the bias is not an artificial feature of the fitting code nor does it appear to be exclusive to the SDSS data. Instead, the bias arises from the simple fact that the majority of the galaxies in our sample are not sufficiently well resolved (their angular sizes

**Table 3.** Best fitting Schechter function parameters for stellar mass functions of total, spheroid and disk components. For the total stellar mass function the best fit parameters are determined using the STY method. For the disk and spheroid stellar mass functions we instead fit a Schechter function to the non-parametric stellar mass function determined using the SWML method—these should be considered useful fitting functions only, not good fits in any statistical sense. For the spheroid and disk stellar mass function fits, the maximum deviation from the SWML data points is given in the final column.

Component	$\log_{10}(M_*/h^{-2}M_\odot)$	$\alpha$	$\phi_0/h^3\text{Mpc}^{-3}$	Max. Dev.
Total	10.92	-1.92	0.00213	N/A
Spheroid	11.03	-0.91	0.00102	50%
Disk	11.06	-1.48	0.00095	40%
SMBH	8.35	-0.97	0.00082	62%

are comparable to the PSF) and so their inclinations are not well constrained. A similar lack of constraint is found in our estimates of the ratio of spheroid to total light, S/T, for galaxies which are poorly resolved. An exception occurs for poorly resolved galaxies which are observed close to edge-on—the inclined disk permits a reasonable separation of disk and spheroid light even in poorly resolved images. We find that excising galaxies with poorly constrained S/T ratios does not significantly affect our conclusion—including these galaxies simply acts to smooth the luminosity function (as we are essentially convolving the true disk/spheroid luminosities with a significant error distribution). A more detailed study of errors on fitted parameters (using the Monte-Carlo techniques described herein, but applied to the full dataset) would permit a better understanding of the effects of these errors to be obtained.

We find that at faint r-band luminosities, the light is dominated by disks whereas at bright luminosities, it is dominated by spheroids, with the changeover occurring at around the characteristic luminosity  $L_*$ . Integrating the luminosity functions, we find the total r-band luminosity densities in spheroids and disks to be  $\rho_L/h = 0.62 \pm 0.08$  and  $\rho_L/h = 1.84 \pm 0.27 L_\odot \text{Mpc}^{-3}$ . Thus, the contribution of the disks to the total luminosity density is three times greater than that of the spheroids. This is in contrast with the findings of previous studies (Schechter & Dressler 1987; Benson et al. 2002), based upon galaxy samples over 40 times smaller, which found the spheroid and disk luminosities to be very nearly equal. Current *a priori* galaxy formation models are able to reproduce these component luminosity functions reasonably well—at least to the extent of predicting the correct trends of abundance with luminosity.

Using an approximate conversion of r-band light to stellar mass, we derive  $\Omega_{\text{stars,disks}} = (0.656 \pm 0.005)h^{-1}10^{-3}$  and  $\Omega_{\text{stars,spheroids}} = (0.383 \pm 0.005)h^{-1}10^{-3}$ , in excellent agreement with the earlier work of Benson et al. (2002) and similar to the fraction of  $55 \pm 2\%$  reported by Tasca & White (2005). From the inferred spheroid mass function and the observed relation between central supermassive black hole mass and spheroid stellar mass, and assuming that all spheroids harbour a central black hole, we infer the supermassive black hole mass function. The associated black hole mass density in the local universe is  $\rho_\bullet = (2.4 \pm 0.13) \times 10^5 h M_\odot \text{Mpc}^{-3}$ , significantly smaller than previous estimates. Improvements in the characterization of the errors in the spheroid mass function and in the stellar

population modelling will enable a better estimate of the spheroid (and black hole) mass function.

We conclude that the local Universe contains around twice as many stars, by mass, in disks than in spheroids. This fundamental ratio is the outcome of the physical processes at play in the formation of the galaxy population.

## ACKNOWLEDGMENTS

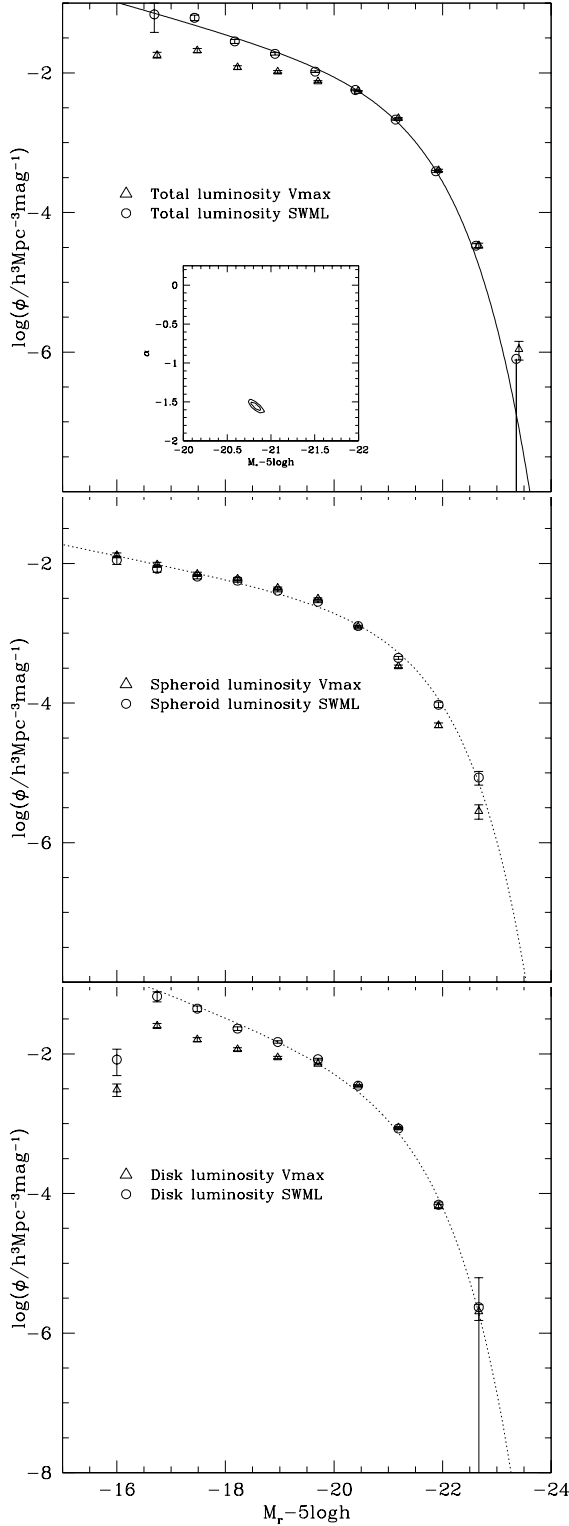
We are extremely grateful to Chris Miller for providing us with all of the SDSS EDR imaging data and Robert Nichol and Tomotsugu Goto for the encouragement given in the early days of this work. We also wish to thank Carlton Baugh for providing his code for computing K+E corrections, Rowena Malbon for providing theoretical predictions for black hole mass functions from the GALFORM semi-analytic model ([www.galform.org](http://www.galform.org)) and Hans Walter Rix and Simon White for many discussions and valuable suggestions. AJB acknowledges support from the Gordon & Betty Moore Foundation, and from a Royal Society University Research Fellowship during part of this work. CSF is the holder of a Royal Society Wolfson Research Merit award.

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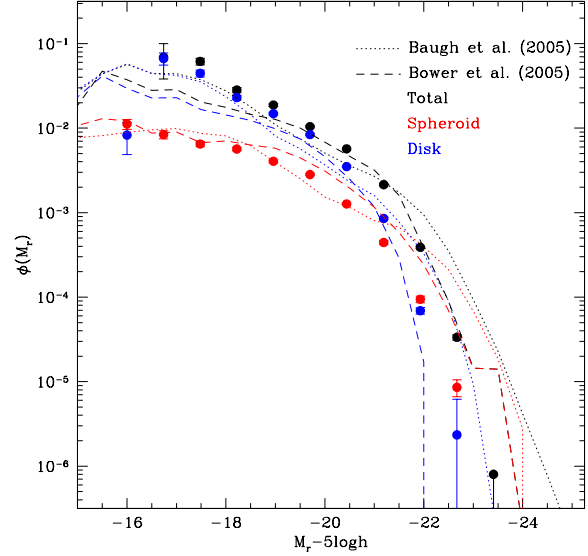
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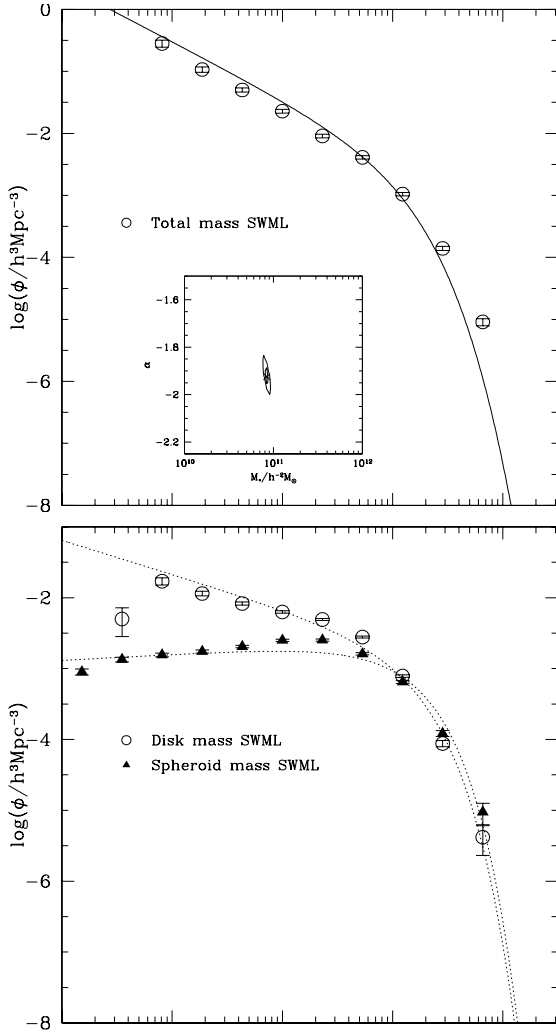


**Figure 14.** Luminosity functions for our sample of 7493 galaxies with  $\chi^2_\nu < 2.0$ . Symbols show the  $V_{\max}$  and SWML estimates; the solid line in the top panel represents the STY fit, while dotted lines in the lower panels indicate the best fit Schechter function to the SWML data points. The top panel displays the total galaxy luminosity function, the middle panel the luminosity function of spheroids and the lower panel, the luminosity function of disks



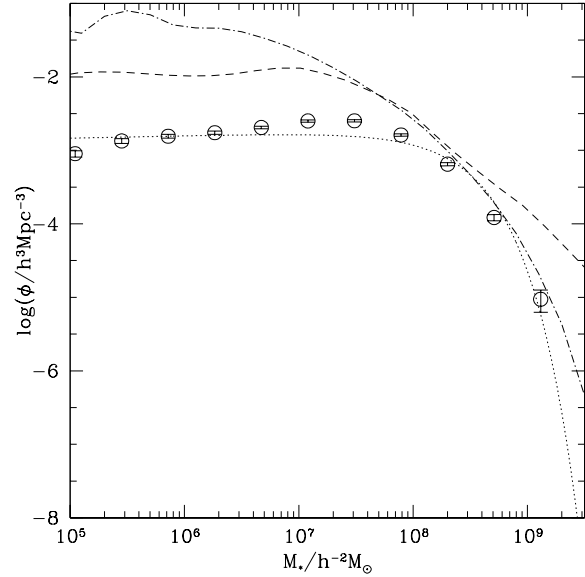
**Figure 15.** Luminosity functions disk and spheroid and total light (indicated by colour; see legend for details). The symbols show our estimates for SDSS galaxies in this work and the lines two different implementations of the GALFORM semi-analytic model: Baugh et al. (2005) (dotted lines) Bower et al. (2006) (dashed lines).

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**Figure 16.** Stellar mass functions for galaxies as a whole and their disk and spheroid components (upper and lower panels) obtained using the SWML estimator. For the total mass function we also plot the Schechter function derived using the STY method (solid line), with the constraints on the parameters  $M_*$  and  $\alpha$  shown in the inset. For the disk and spheroid mass functions dotted lines shows the Schechter function which best fits the SWML data points.

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**Figure 17.** The mass function of supermassive black holes in galactic spheroids. Symbols show the black hole mass function implied by our observationally determined spheroid stellar mass function assuming that  $M_*/M_\odot = 1.6 \times 10^8 [M_{\text{spheroid}}/10^{11} M_\odot]^{1.12}$  (Häring & Rix 2004). The dotted line shows the Schechter function which best fits the SWML data points. Other lines show results from the galaxy formation model of Malbon et al. (2006) when using the parameters specified by Baugh et al. (2005; dot-dashed line) and Bower et al. (2006; dashed line).

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## APPENDIX A: METHODS FOR QUANTITATIVE GALAXY MORPHOLOGY

In this Appendix, two independent methods for spheroid-to-disk decompositions of galaxies are described and compared. A number of tests are performed to reveal and estimate any potential biases in the decomposition codes.

### A1 Introduction to spheroid-to-disk decomposition

To first approximation, the surface brightness of a galaxy can be expressed as the sum of a highly concentrated central component, the spheroid, and an extended disk. Empirically, the surface brightness profiles of spheroids and disks are well represented by the following functions,

$$I_s = I_e \exp(-7.67[(r/r_e)^{1/4} - 1]), \quad (\text{A1})$$

an  $r^{1/4}$ -law, for the spheroid (de Vaucouleurs 1961), where  $r_e$  is the half-light radius and  $I_e$  is the surface brightness at  $r_e$ , and,

$$I_d = I_0 \exp(-r/r_d), \quad (\text{A2})$$

an exponential-law for a face-on disk, where  $r_d$  is the exponential disk scale-length and  $I_0$  is the central intensity.

Equations A1 and A2 can be used to construct model images of the galaxy. Comparison with the surface brightness distribution of each galaxy, including the effects of inclination, enables the fitting parameters to be determined.

Andredakis et al. (1995) used this technique to fit the spheroid components of a sample of morphologically selected galaxies with types ranging from S0 to Sbc. They assumed a more general type of profile,

$$I_b = I_e \exp(-b_n[(r/r_e)^{1/n} - 1]), \quad (\text{A3})$$

first proposed by Sérsic (1968), where  $n$  is often referred to as the Sérsic index and determines the ‘peakiness’ of the profile and  $b_n$  is a constant dependent on the value of  $n$ . Andredakis et al. found that the value of  $n$  varied systematically from 1 for late-type spheroids to 6 for early-type spheroids. de Jong (1996) also suggested that the spheroids of field spirals are better fit using a pure exponential (i.e.  $n = 1$ ) profile.

### A2 Methods for 2-dimensional spheroid-to-disk decomposition

#### A2.1 Fitting parameters revisited

On a Cartesian grid  $(x, y)$  the effective  $r$  in eqns. (A1) and (A2) become

$$r^2(x, y) = \frac{1}{e_s} [x \cos(\theta_s) - y \sin(\theta_s)]^2 + e_s [x \sin(\theta_s) + y \cos(\theta_s)]^2 \quad (\text{A4})$$

$$r^2(x, y) = [x \cos(\theta_d) - y \sin(\theta_d)]^2 + \frac{1}{\cos(i)^2} [x \sin(\theta_d) + y \cos(\theta_d)]^2 \quad (\text{A5})$$

and

$$r^2(x, y) = [x \cos(\theta_d) - y \sin(\theta_d)]^2 + \frac{1}{\cos(i)^2} [x \sin(\theta_d) + y \cos(\theta_d)]^2 \quad (\text{A6})$$

respectively. Here  $\theta_s$  and  $\theta_d$  are the spheroid and disk position angles, where a position angle is defined as the angle of orientation of the galaxy’s main axis with respect to some coordinate system, and  $e_s$  is the spheroid ellipticity used to describe the deviation from circularity of the spheroid component

In terms of a  $r^{1/n}$  spheroid and an exponential disk, and including the sky background, 2D decomposition usually requires a total of 13 free parameters:

- total flux in the galaxy;
- S/T: ratio of the amount of light in the spheroid to the total amount of light;
- $r_e$ : effective radius of the spheroid;
- $e_s$ : spheroid ellipticity;
- $\theta_s$ : spheroid position angle;
- $r_d$ : scale length of the disk;
- $i$ : inclination angle of the disk;
- $\theta_d$ : disk position angle;
- $x_c, y_c$ : subpixel offset of the galaxy centre;
- residual sky background level;
- FWHM of the point spread function (PSF);
- $n$ : Sérsic index.

In order to make the decomposition procedure as accurate and fast as possible, the following points must be taken into account:

- The PSF smooths the galaxy image and so to achieve accurate fits, it must be modelled accurately and included in the mock images.
- The fit is carried out using small “thumbnail” regions around each galaxy. The size of the thumbnail must be small enough to enable a fast fit to the image, but large enough to include all regions of the galaxy with significant signal-to-noise.



- The mean sky background level should be  $\sim 0$  since the decomposition codes are designed to work with no (or very little) background; any excess sky light can be mistaken for galaxy light and lead to incorrect parameter estimation.

We now explore the similarities and differences of two independent multi-dimensional fitting codes, the GIM2D code of Simard et al. (2002) and the GALACTICA code of Benson et al. (2002).

### A2.2 Galaxy image 2D decomposition (Gim2D)

GIM2D is a publicly available code written by Simard et al. (2002) and widely used for automated spheroid-to-disk decompositions of galaxy light profiles (Balogh et al. 2002; Nelson et al. 2002). This code was purposely written for imaging with the Hubble Space Telescope Wide Field and Planetary Camera which has a very well modelled PSF (Krist 1995). The code can also be used for ground-based imaging data but, in this case, special attention must be given to the much larger and less well defined PSF.

#### Object Detection

To extract a postage-stamp image around each galaxy, GIM2D relies upon SExtractor (Bertin & Arnouts 1996). SExtractor determines the galaxy centroid position and the area at the faintest detected isophote to be obtained and also measures the mean level of the sky background for each galaxy (a  $3\sigma$  threshold is usually sufficient to discriminate between the object and the background). GIM2D extracts a postage-stamp of size equal to a multiple of a galaxy isophotal area. A value of  $15 \times \text{iso\_area}$  was found to be optimum. The sky-background is not recommended to be treated as a free fitting parameter in GIM2D because the underlying sky is not well known and can potentially bias the output (Simard et al. 2002). However, before the decomposition procedure is initiated, GIM2D uses the pixels flagged by SExtractor as belonging to the background (flag value 0) to recompute the background value, therefore ensuring that the mean sky level is close to zero. All the background pixels and also pixels flagged as ‘bad’ (flag value -2) by SExtractor are subsequently excluded from the fitting altogether.

#### Point Spread Function

During the minimization in GIM2D, the PSF is kept fixed. For ground-based imaging, the PSF is obtained from a bright unsaturated stellar image; for HST data, an analytic PSF modelled using the Tiny Tim software is used (Krist 1995).

#### Minimization Technique: Metropolis Algorithm

GIM2D allows up to 12 parameters to be fit<sup>4</sup> and uses the Metropolis algorithm (Metropolis et al. 1953) to search for the minimum  $\chi^2$  in this multi-dimensional parameter space. Before starting the Metropolis algorithm, GIM2D works in the Initial Condition Filter (ICF) mode, i.e. it creates a user-specified number of models between the limits of the user-specified multi-dimensional parameter space. The

ICF computes the model likelihoods and sets the sampling origin to the parameters of the best model, making it a sub-volume to be exploited by the Metropolis Algorithm.

#### GIM2D Outputs

After finding the model that corresponds to the highest likelihood, GIM2D produces a residual (object - model) map and calculates the value of the corresponding  $\chi^2_\nu$ . If  $\chi^2_\nu \sim 1$  and the residual map has no remaining galaxy structure, the best-fit model is accepted.

### A2.3 GALACTICA

#### Introduction

The 2-dimensional decomposition code described here is based on a technique proposed by Wadadekar et al. (1999). GALACTICA was developed by Benson et al. (2002) and assumes the standard empirical formalisms for the 2-dimensional surface brightnesses of a galaxy spheroid and disk components respectively (eqns. A1 & A2).

#### Object detection

To locate and extract a postage-stamp image around every galaxy GALACTICA employs the same method as GIM2D. SExtractor is also used to measure the mean level of the sky background for each galaxy (a  $3\sigma$  threshold is usually sufficient to discriminate between the object and the background) and this value is subtracted from the corresponding galaxy image. To mask any overlapping objects within the extracted postage-stamp GALACTICA relies upon an in-built masking algorithm which finds any objects that contaminate the galaxy of interest and masks them out. The galaxy itself is also detected by the algorithm using a  $5\sigma$  threshold above the sky background. Pixels which have not been flagged as belonging to any of the detected objects are used in the sky background fitting.

#### Point spread function: Moffat profile

To correct for the effect of seeing the GALACTICA code generates a Moffat profile star image (Moffat 1969) of a given full-width half-maximum (FWHM) expressed in terms of  $\sigma_{\text{PSF}} = \text{PSF}_{\text{FWHM}}/2.35$ . This analytic profile is defined by

$$\text{PSF}(r) = \text{const}/[1 + (r/\alpha)^2]^\beta \quad (\text{A7})$$

and is thought to represent the overall PSF shape better than a pure Gaussian which only approximates the core regions. Here  $\alpha$  represents the width of the PSF and is related to the  $\text{FWHM} = 2\alpha\sqrt{2^{1/\beta} - 1}$  (Trujillo et al. 2001).  $\beta$  governs how ‘peaky’ the PSF profile is (the larger  $\beta$  is, the more Gaussian-like the profile becomes). A value of  $\beta = 4.5$  is used throughout. The  $\alpha$  parameter can be fine-tuned to a particular dataset using the average FWHM for the data. GALACTICA lets  $\sigma_{\text{PSF}}$  be a free fitting parameter to allow for any small changes in the PSF between the position of the stars in the image and the galaxy position.

#### Minimization technique: Powell’s method

The code requires explicit initialization of the fitting parameters. The initial value of the S/T ratio is always fixed at 0.5, although starting with a value of S/T randomly distributed in the range 0 to 1 has no effect on the recovered S/T distribution. The starting position angles of the disk

<sup>4</sup> Sérsic index is held fixed at 4.

and spheroid components, their characteristic radii and the disk inclination angle are calculated directly from the image. The ranges over which parameters are allowed to vary during the fitting are specified and fitting outside these limits is not possible.

$\chi^2$  is minimized in a 12-parameter space; the Sérsic index is set to  $n = 4$  and the FWHM of the PSF and the residual sky background level are additional fitting parameters not included in GIM2D. The minimization routine is also rather different from the one used by GIM2D. In GIM2D every parameter is varied at each step according to the ‘temperature’ of the fit. In GALACTICA one parameter is minimized at a time, i.e. all but one parameters are ‘frozen’ until a minimum for this parameter is found and the process is repeated for the entire set of parameters until the global minimum is found —the essence of Powell’s method (see Press et al. 1992 for further details). This method is good for finding a global extremum but is typically slower than the Metropolis algorithm employed by GIM2D.

#### GALACTICA outputs and error estimation

After convergence is achieved, the best-fit parameters are output along with the best-fit model image and the residual map obtained by subtracting the model galaxy from the real image. The value of  $\chi^2$  per degree of freedom,  $\chi^2_\nu$ , is then calculated. Errors on the fitted parameters are obtained using a Monte Carlo method: we create 30 realisations of the best fitting model for each galaxy by adding random noise, drawn from a Poisson distribution, to the model image. The distribution of the best-fitting parameters of the model realizations is then used to estimate the uncertainty in the fit. This method allows the uncertainties in the image parameters to be obtained without any assumptions about their distribution.

### A3 GIM2D vs. GALACTICA comparison

#### A3.1 Introduction

The previous section described two independent codes for estimating basic galaxy structural parameters. Although these codes assume the same analytic surface brightness profiles to fit the spheroid and disk components, the differences in the number of fitting parameters and in the minimization techniques are sufficiently large to make a comparison interesting and important.

We first note the following points relevant to the comparison:

- both codes assume a fixed value of  $n = 4$ ;
- the sky background is always kept fixed by GIM2D although the code is allowed to recompute and correct the background level before the minimization procedure starts;
- GALACTICA always treats the sky background as a free parameter;
- the ellipticity is defined differently: GIM2D fits  $e = 1 - b/a$  while GALACTICA fits  $a/b$ ;
- the seeing is fixed in GIM2D but in GALACTICA it is allowed to fluctuate between  $\pm 5\%$  of the specified  $\sigma_{\text{PSF}}$ ;
- the position angles in GIM2D are defined with respect to the  $y$ -axis of a Cartesian system while in GALACTICA they

**Table A1.** Parameter ranges used for constructing mock galaxies.

Parameter	Low limit	High limit
S/T	0.0	1.0
$r_{\text{e,d}}$ (pixels)	1	12
$e_{\text{s}}$	0.0	0.8
$i$ (degrees)	0.0	90.0
$\theta_{\text{s,d}}$ (degrees)	0.0	180.0
FWHM (")	1.4	1.4

are defined clockwise from the  $x$ -axis. (The position angles of the spheroid and disk are allowed to vary in both codes; a large difference between these can be a signature of a barred structure; Simard et al. 2002).

To quantify the performance of the codes, a series of tests were conducted as we now describe.

#### A3.2 Tests using model galaxies

A useful in-built feature of both GALACTICA and GIM2D is the ability to create model galaxies. The initial tests and code comparisons described below were carried out on model galaxies generated “internally” by GALACTICA.

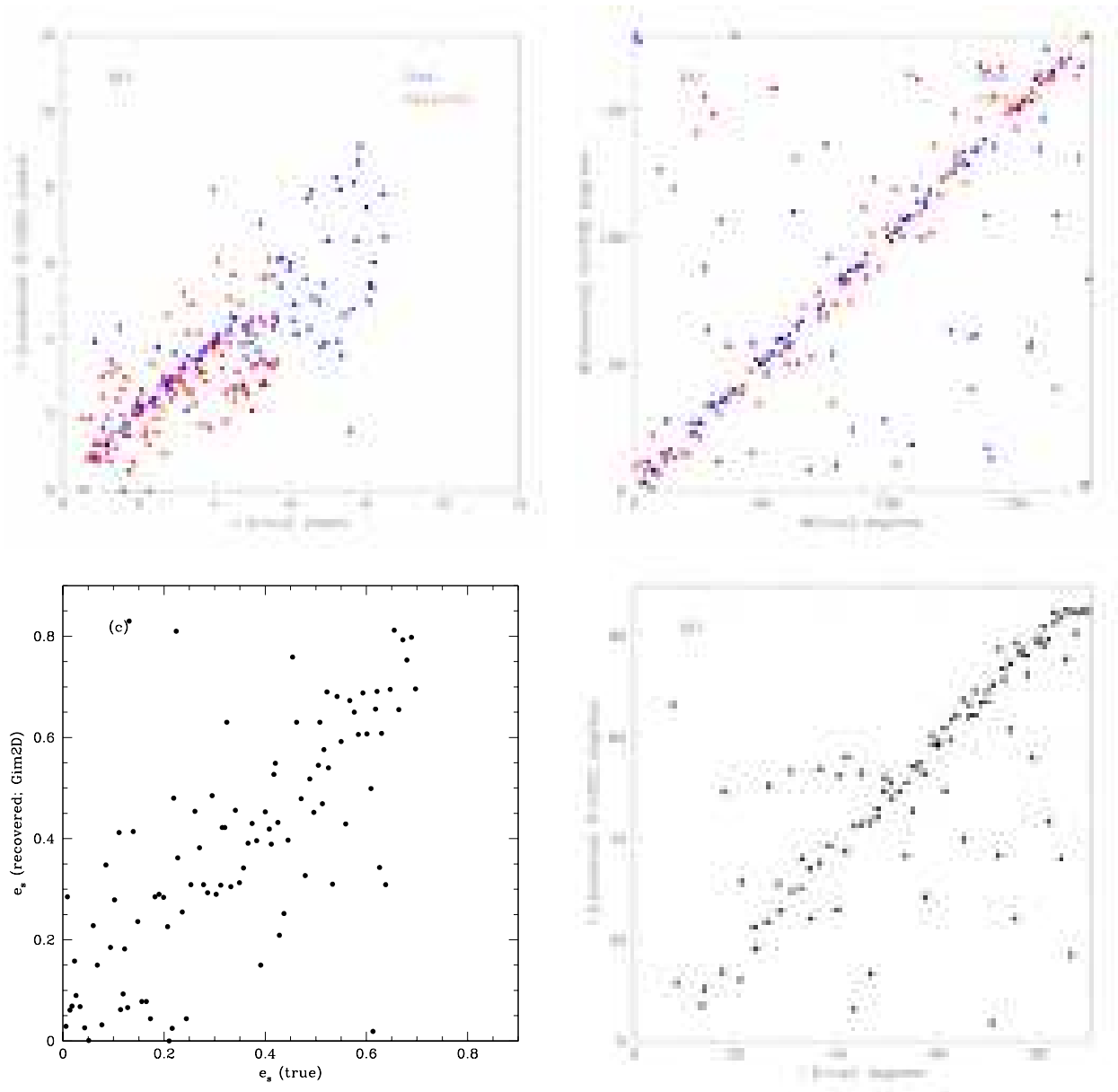
Model galaxies were constructed adopting parameter values chosen at random between realistic limits (Table A1) and matching the total counts measured in a typical real galaxy. Poisson noise was added to the model galaxy after its image was convolved with an analytic Moffat PSF corresponding to a typical value of the seeing. This PSF is subsequently used as the GIM2D PSF. Model galaxies were analyzed with both codes using exactly the same procedures as for real galaxies.

Comparison of the known input S/T values and the values output by GIM2D and GALACTICA for 100 model galaxies are shown in Figs. A1 and A2 respectively. In both cases the codes recover the input S/T very well. The scatter in the recovered S/T ratios is  $\sigma_{\text{rms}} \sim 0.10$ . The remaining parameter recoveries are shown in Figs. A3 and A4.

The results shown in Figs. A1 to A4 demonstrate that both GIM2D and GALACTICA produce, on the whole, reliable spheroid/disk decompositions for a set of artificial galaxies generated according to the model assumed by the code. A more stringent test of the codes is to apply them to model galaxies generated independently of the codes themselves. Model galaxies were therefore externally created using the IRAF task MKOBJ.

The model parameters were taken from Table A1; a model galaxy is created for given values of the size, orientation and ellipticity (in this case defined as  $b/a$ ) and the image is convolved with a specified seeing. A useful feature of this approach is that a real science frame can be fed into MKOBJ and the model galaxy added to a blank patch of the sky on this science frame. By matching, on average, the counts in a real galaxy, this procedure ensures that the artificial image has the similar noise characteristics to the real data.

To generate the model galaxies several SDSS  $r'$ -band frames were extracted, each typically containing  $\sim 5$  SDSS catalogued galaxies. Each frame was taken from a different patch of the sky. Counts associated with the SDSS galaxies

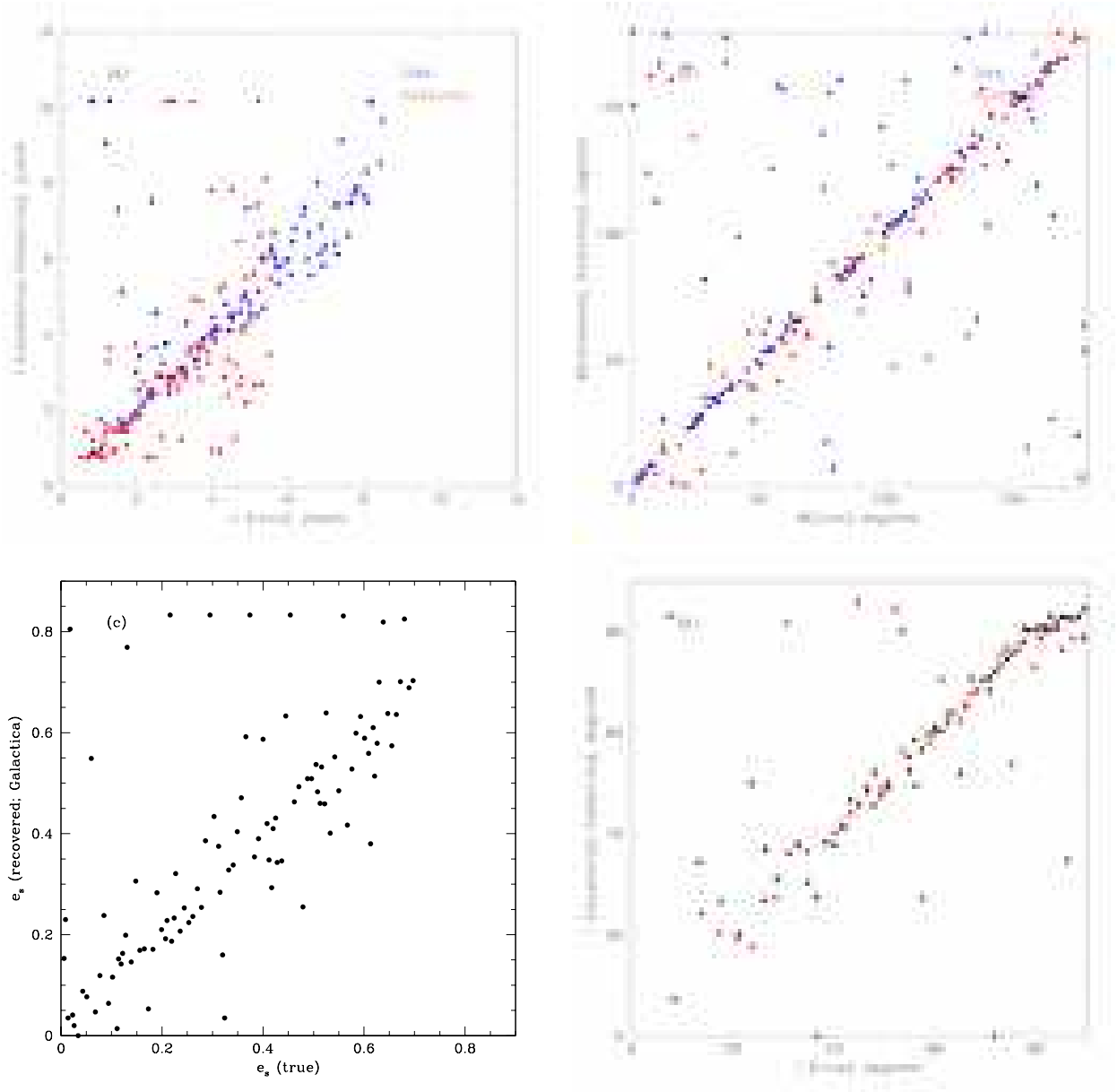


**Figure A3.** Correlations between input and recovered values of the characteristic radius, position angle, ellipticity and inclination for 100 model galaxies created using the GALACTICA code and decomposed using GIM2D. a) Disk and spheroid radii. b) Disk and spheroid position angles. c) Spheroid ellipticity. d) Disk inclination. There is a saturation at  $i = 85^\circ$  which is the upper limit that GIM2D allows for the disk inclination. All other parameters correlate well although a significant scatter is seen.

were measured using the SEXTRACTOR `flux_best` parameter. In order to test the decomposition algorithm over a realistic range of galaxy properties, this procedure was applied to galaxies spanning a range of apparent magnitude and apparent shape and size using the IRAF task MKOBJ, the model galaxies were inserted across the blank regions of the sky in the original frames. The postage-stamps for these galaxies were extracted from both the science and SEXTRACTOR frames and the decomposition codes run treating the extracted model galaxies just as the real ones.

The results of the GIM2D decompositions of the model galaxies are shown in the Fig A5. The agreement between

the input and the output S/T ratios for the pure exponential disks ( $S/T = 0$ ) is excellent. However, the recovered S/T ratio for  $S/T = 0.5$  is biased by  $\Delta S/T = 0.1$  and for  $S/T = 1.0$  it is biased by  $\Delta S/T = 0.2$ . The tendency is always to underestimate the amount of spheroid or, equivalently, overestimate the amount of disk. GIM2D can be fine-tuned to recover the input S/T ratios with  $\Delta S/T \simeq 0.1$  across the full S/T range. For this, GIM2D requires that the size of the zone around the lowest SEXTRACTOR isophote used in the re-calibration of the sky background should be set to  $\sim 30$  pixels (the default value is 10 pixels). This ensures that any faint galaxy flux does not contribute to the re-calibrated background flux thus minimizing any bias in S/T.

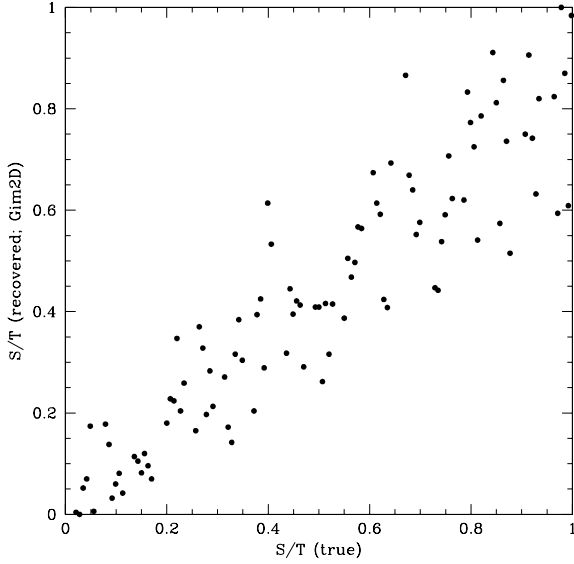


**Figure A4.** Correlations between input and recovered values of the characteristic radius, position angle, ellipticity and inclination for 100 model galaxies created and decomposed using the GALACTICA code. a) Disk and spheroid radii. Note that the reconstruction seems to hit the upper limit on the characteristic radii when the input radii are very small or very large. b) Disk and spheroid position angles. c) Spheroid ellipticity. Note that a number of very elliptical spheroids are found. These are mostly for galaxies which have very small spheroids. d) Disk inclination. Note an apparent saturation at  $i \sim 85^\circ$ . Even though the GALACTICA code allows the disk to be fully inclined ( $i = 90^\circ$ ), the reconstruction avoids this upper limit.

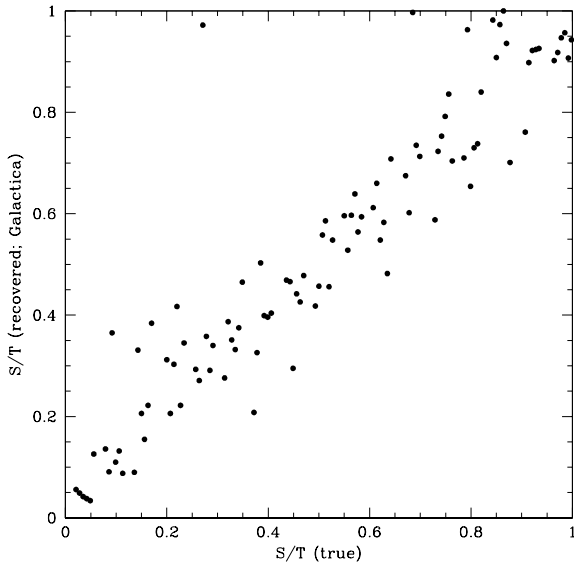
The results of the GALACTICA decompositions of the model galaxies are shown in Fig. A6. For the pure exponential disks the recovered S/T ratios are, again, very good. As before for larger S/T, this ratio is overestimated and peaks at S/T = 1.0, implying that many pure  $r^{1/4}$  galaxies have acquired a fictitious disk component. There appears to be no correlation between the size of the overestimation and the magnitude or scale radii of the input galaxy but a weak correlation with the minor/major axis ratios: the S/T deviation is largest for the most elliptical profiles. The most prominent correlation, however, is between the recovered S/T ratio and

the sky background. As discussed earlier, the GALACTICA code allows the sky background to fluctuate a little to allow for uncertainties in the estimated background. The fact that the deviation between the input and the output S/T ratios is largest when the ‘fitted’ background is smallest implies that an extraneous disk component is found where, in fact, the extra counts are due to the sky background. Since there is no sharp cut-off in either the spheroid or the disk, at the galaxy edges the galaxy surface brightness profile and the sky background are indistinguishable.

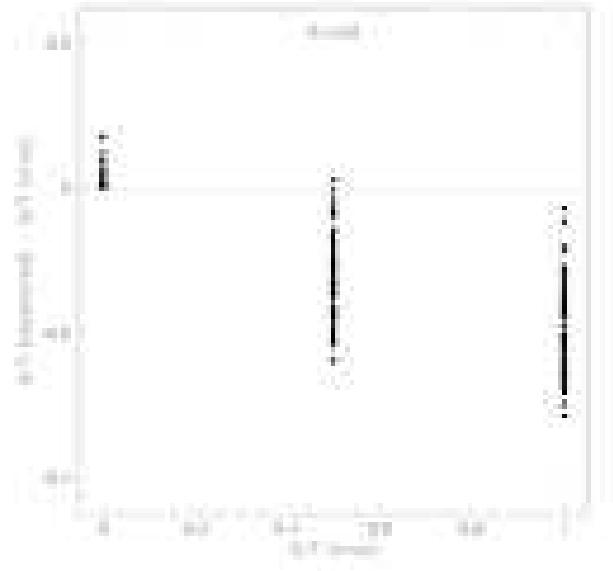
GIM2D performs marginally better than GALACTICA in



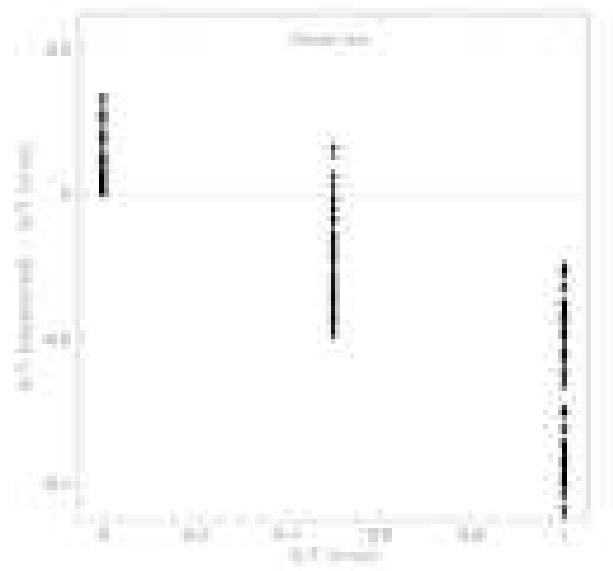
**Figure A1.** Correlation between the input S/T ratios for 100 model galaxies and the best-fit S/T ratios recovered using GIM2D. The scatter around the mean is  $\sigma_{\text{rms}} = 0.10$ . Note a small systematic underestimate of S/T, accompanied by increased scatter, for large input values. The recovered characteristic radii of these galaxies are the largest and the sky background recalculated by GIM2D for these galaxies is high. The resulting confusion between the background and an extended surface brightness profile accounts for this effect.



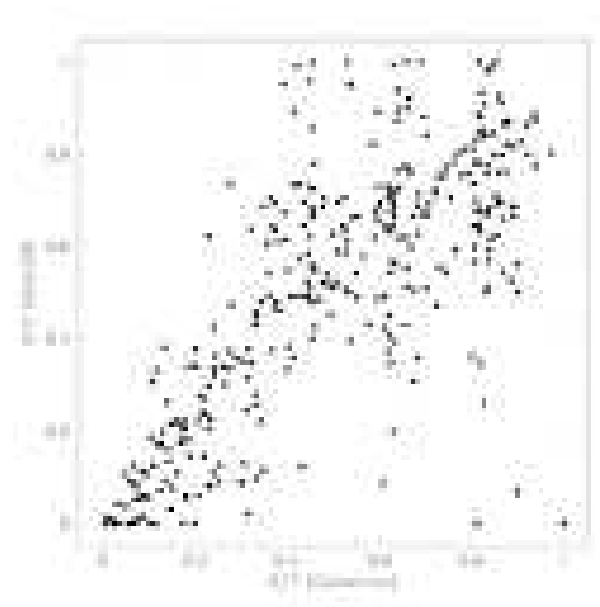
**Figure A2.** Correlation between the input S/T ratios for 100 model galaxies and the best-fit S/T ratios recovered using GALACTICA. The scatter around the mean is  $\sigma_{\text{rms}} = 0.11$ .



**Figure A5.** The difference between input and recovered values of S/T using GIM2D for a set of 250 model galaxies. The mean offset in the recovered values are for  $\Delta S/T = 0.02$  (for  $S/T = 0.0$ ),  $\Delta S/T = 0.13$  (for  $S/T = 0.5$ ) and  $\Delta S/T = 0.20$  (for  $S/T = 1.0$ ). Model galaxies span a range of apparent magnitudes, sizes and orientations.



**Figure A6.** The difference between input and recovered values of S/T using GALACTICA for a set of 250 model galaxies. The mean offset in the recovered values are for  $\Delta S/T = 0.05$  (for  $S/T = 0.0$ ),  $\Delta S/T = 0.11$  (for  $S/T = 0.5$ ) and  $\Delta S/T = 0.24$  (for  $S/T = 1.0$ ). Model galaxies span a range of apparent magnitudes, sizes and orientations.



**Figure A7.** The correlation between the S/T ratios of a sample of 350 SDSS galaxies inferred using GIM2D and GALACTICA. The Spearman rank correlation coefficient of 0.74 indicates a significant correlation. Note that almost no pure disks (i.e. S/T= 0.0) are detected by GALACTICA but a few are detected by GIM2D. Most of these galaxies have spheroid characteristic radii of less than 2 pixels (as does the single GALACTICA detection at S/T= 1.0). GIM2D finds larger characteristic radii for these galaxies. The overall scatter is  $\sigma_{\text{rms}} = 0.19$ , but there is a marked increase in the scatter for larger S/T ratios.

the recovery of the S/T ratios. However, both codes show similar biases in the decompositions. Their relative performance on a set of real galaxies is compared next. This allows a more realistic comparison of the codes but, of course, there is no *a priori* correct answer.

### A3.3 Tests using real galaxies

To ensure a uniform sampling of the [S/T, apparent magnitude] space, the comparison was carried out using a subsample of SDSS galaxies selected in bins of 0.5 in apparent magnitude and 0.2 in S/T ratio (as determined by GALACTICA). Unsaturated stellar images with high S/N were extracted from the SDSS galaxy frames and used in the GIM2D PSF deconvolution. The GALACTICA Moffat PSF was fine-tuned to fit the SDSS data well. Fig. A7 demonstrates a significant correlation (Spearman rank correlation coefficient of 0.74) between the S/T ratios for  $\sim 350$  SDSS galaxies inferred using GIM2D and GALACTICA. There are no systematic differences between the results of the two codes. The correlations between other parameters are displayed in Fig. A8.

## APPENDIX B: GALACTICA METHODOLOGY

### B1 Code speed limitations and galaxy binning

The SDSS galaxy sample contains less than 200 galaxies whose postage-stamp size exceeds 91 pixels on a side. To reduce the processing time, these large postage-stamps were binned  $2 \times 2$ . This has the added advantage that large nearby bright galaxies are sampled at a resolution comparable to that of more distant objects. To make sure that the binning procedure does not bias the recovery of the galaxy S/T ratios, we carried out a series of tests.

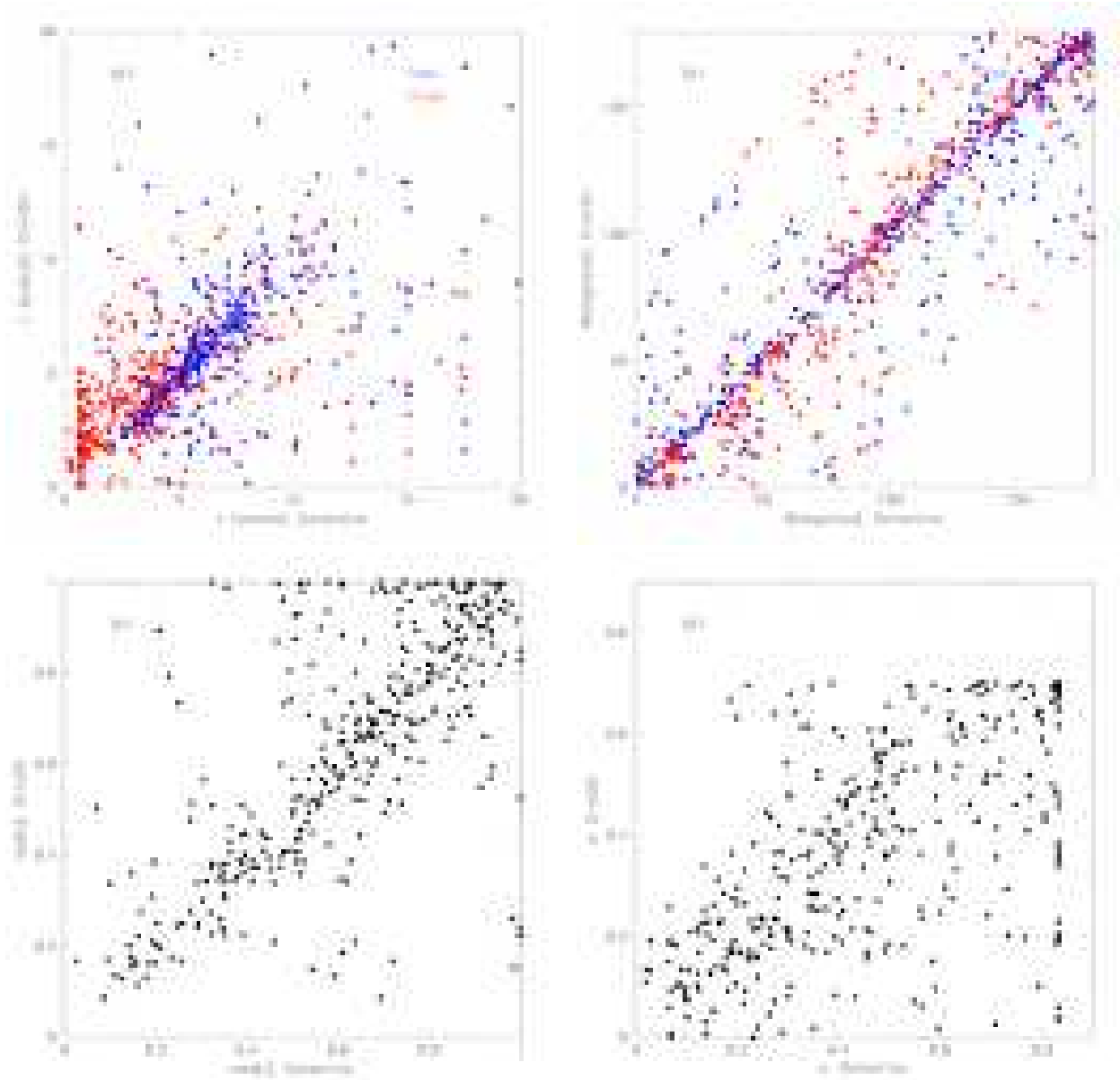
Several model galaxies were created using our standard procedure (see Appendix A3.2) and binned using the IRAF task `BLKAVG`. GALACTICA was then used to perform the fitting ensuring that the pixel for the binned image is set to  $2 \times$  the normal pixel size ( $2 \times 0.396''$ ) and that the noise properties in this ‘super-pixel’ are changed accordingly. Figs. B1 and B2 show fits to model galaxies consisting either of a pure exponential disk or a pure  $r^{1/4}$  spheroid after the original model images were binned by  $2 \times 2$ . In the case of the pure exponential disk, the fit to the model is perfect. The fit to the pure  $r^{1/4}$  galaxy, however, shows a similar bias to that seen for the unbinned data (Appendix A) as indicated by the low value of the recovered S/T= 0.8. This shows that the binning in itself is not responsible for the observed S/T bias. The  $2 \times 2$  binning was therefore applied to all the SDSS galaxies whose postage-stamps are greater than  $91 \times 91$  pixels.

The binning works very well if the binned galaxy does not exhibit much internal structure (as in the model galaxies). However, for a galaxy which exhibits significant internal structure, a fit with  $\chi^2_\nu > 2.0$  is more typical. Whether decomposing such galaxies even without binning would lead to a good fit is unclear as demonstrated for 2 SDSS galaxies in Fig. B3 and B4. The top images in both figures show the postage stamp and the residual map for the unbinned galaxy which has size  $101 \times 101$  pixels. (To speed up the calculation, the postage stamp was trimmed by 5 pixels on either side.) The bottom panels show the corresponding images for the binned versions of the same galaxy.

The galaxy in Fig. B3 exhibits much more internal structure than the galaxy in Fig. B4 as is clearly from both the value of  $\chi^2_\nu$  and the residual image. This supports the conclusion that galaxies which exhibit internal structure are poorly fit irrespective of whether they are binned or not. The recovered S/T ratios for the unbinned and binned data, although different, are consistent with the typical errors in the S/T ratio. We conclude that the S/T distribution of the final SDSS galaxy sample is not biased by binning these large, bright nearby objects (most of which contribute to the faint-end of the luminosity function – see Section 4).

### B2 SDSS data and the goodness-of-fit

The selected SDSS sample of 8839 galaxies is too large for each of the residual images to be inspected by eye to ensure a satisfactory decomposition as suggested by the  $\chi^2_\nu < 2.0$ . However, a randomly selected sample of residuals was examined by eye to ensure that they were indeed predominately noise dominated. The  $\chi^2_\nu < 2.0$  criterion was therefore adopted to define a ‘well fit’ dataset of 7493 galaxies. To



**Figure A8.** Correlation between various parameters inferred using GIM2D and GALACTICA for a sample of 350 SDSS galaxies. No systematic differences in the recovered parameters are apparent although the scatter can be quite large.

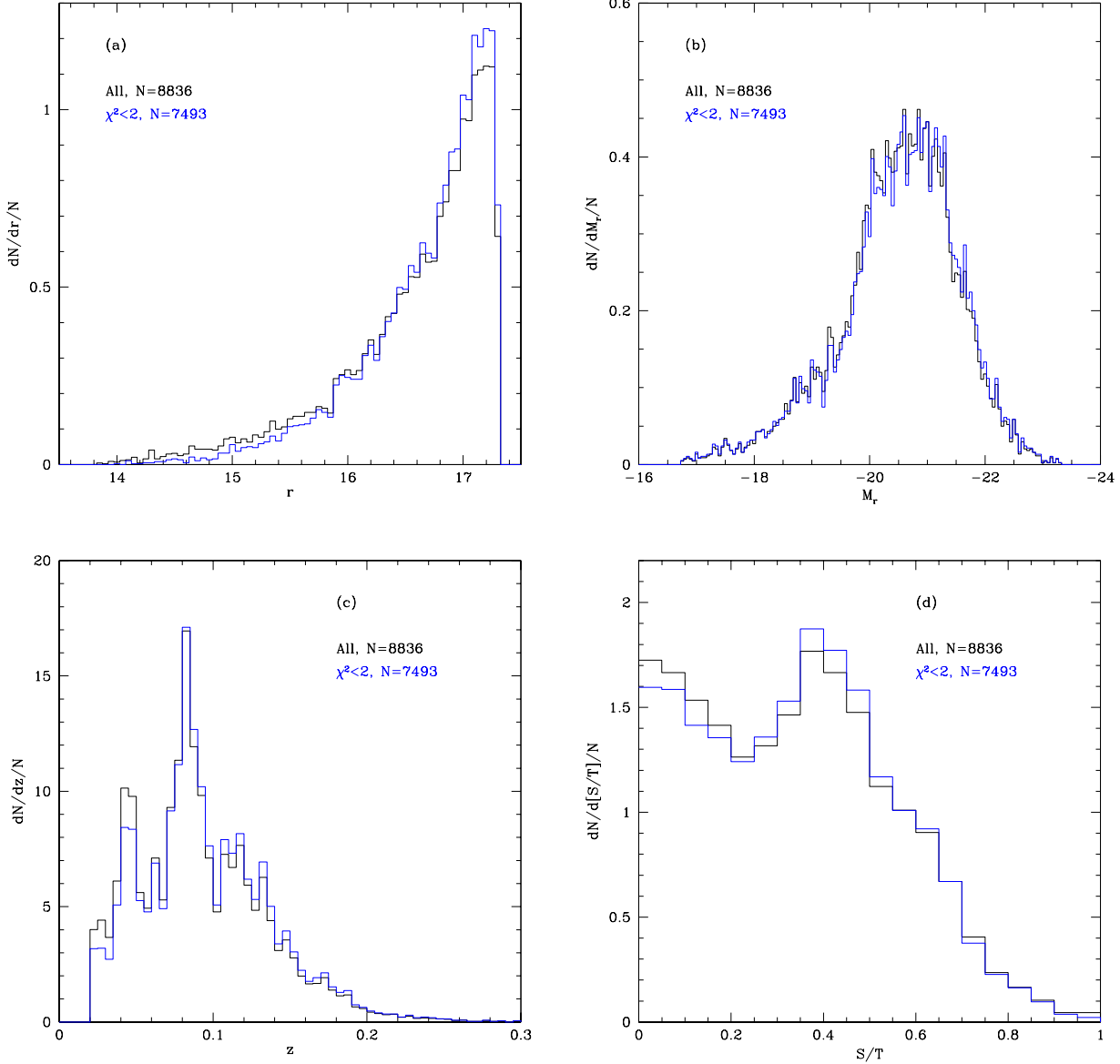
check that no selection biases were introduced by the rejection of galaxies with  $\chi^2_\nu > 2.0$ , we compare the distributions of some basic properties of these galaxies and of the well-fit subset in Fig. B5. The figure shows a deficit of objects with  $S/T > 0.7$ . This is most likely due to the bias in the GALACTICA code discussed in Appendix A. The significance of this bias and its influence on the final results is discussed in the main body of the paper.

### B3 S/T error estimates

Benson et al. (2002) developed a Monte Carlo approach to estimate the errors on the fitted parameters in GALACTICA. This method is very time consuming, requiring several CPU days for a typical galaxy. A full Monte Carlo analysis is

therefore impractical for a large dataset such as the one in this paper. Instead, to obtain representative error estimates, we split the sample into bins of apparent magnitude of width 0.5, and selected from each of these bins 5 galaxies from each of three further bins in  $S/T$  ( $0.0 < S/T < 0.3$ ,  $0.3 < S/T < 0.6$ ,  $0.6 < S/T < 1.0$ ). The full Monte Carlo analysis was performed on the selected subsample and the medians of the derived errors taken to be representative for galaxies in each  $[r_{\text{mag}}, S/T]$  bin.





**Figure B5.** Distributions of (a) apparent magnitude, (b) absolute magnitude, (c) redshift and (d) recovered S/T ratios. Results for the total sample are shown by the black histograms and for the sample with  $\chi^2_{\nu} < 2.0$  by the blue histograms. The two distributions are similar, indicating that excluding poorly fit galaxies does not introduce any obvious biases in the sample.

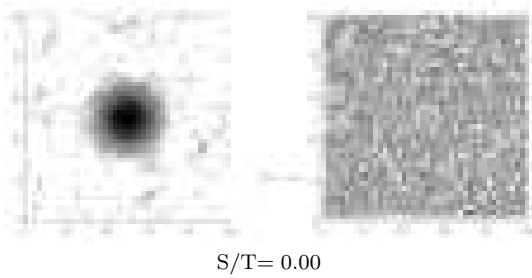
## APPENDIX C: SPHEROID-TO-DISK RATIOS AND GALAXY MORPHOLOGIES

### C1 Morphological classification using colour

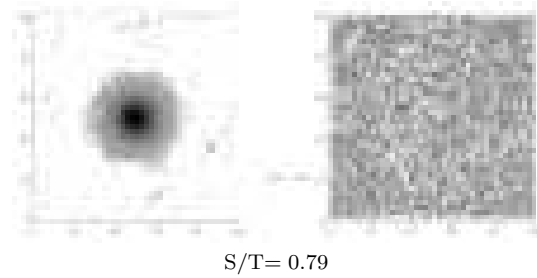
It has long been known that galaxy colour is a useful indicator of whether a galaxy is elliptical (old, red) or spiral (young, blue) (de Vaucouleurs 1961) since the dominant stellar populations are reflected in the galaxy colours. Investigating the colour-magnitude and colour-colour diagrams, Strateva et al. (2001) have shown that the  $(u-r)$  colour distribution of SDSS galaxies has two maxima which are separated by a well-defined minimum at  $(u-r) = 2.2$  and that 98% of galaxies spectroscopically classified as ‘early’ types have  $(u-r) > 2.2$  whilst 73% of spectroscopically classified

‘late’ types have  $(u-r) < 2.2$ . Strateva et al. (2001) have also shown that this separator also applies for a subsample of visually classified morphological types where 80% of galaxies visually classified as E, S0 or Sa have colours redder than  $(u-r) = 2.2$  and 66% of galaxies visually classified as Sb, Sc and Irr have colours bluer than  $(u-r) = 2.2$ . The  $(u-r)$  separator has already been used to study morphological properties of galaxies in the SDSS sample as a function of environment (Goto et al. 2002; Balogh et al. 2004).

Our derived S/T ratios are plotted against  $u-r$  colour in Fig. C1. The B/T distributions of red galaxies and blue galaxies are significantly displaced relative to each other: most galaxies classified by colour as early types ( $u-r > 2.2$ ) have  $S/T > 0.4$  whereas most galaxies classified by colour



**Figure B1.** Fit to a pure exponential model galaxy after the original  $101 \times 101$  model image was binned  $2 \times 2$ . The recovered  $S/T = 0$ , corresponding to a pure exponential. The good fit is evident from both the  $\chi^2_\nu \sim 1$  and the noise-dominated residual image.

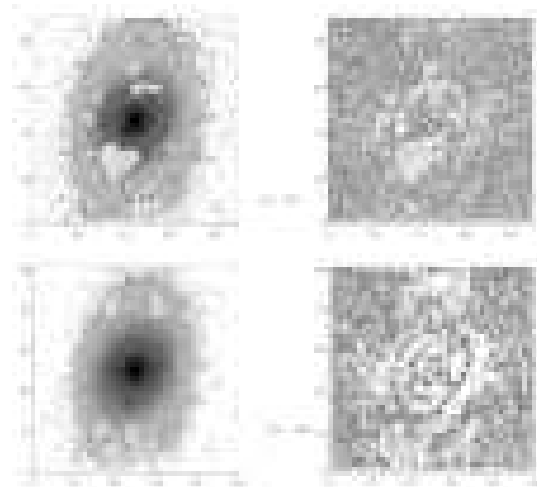


**Figure B2.** Fit to a pure  $r^{1/4}$  model galaxy after the original  $101 \times 101$  model image was binned  $2 \times 2$ . The recovered  $S/T = 0.8$ , showing once more that the GALACTICA code returns a biased estimate of  $S/T$ . The good fit is inferred from both  $\chi^2_\nu \sim 1$  and the noise-dominated residual image.

as late types ( $u - r < 2.2$ ) have  $S/T < 0.2$ . However, the red galaxies in particular span a large range in  $S/T$  ratio. This suggests that blue star-forming galaxies are predominantly disk-dominated but that disk-dominated galaxies include both star-forming (blue) and passive (red) galaxies<sup>5</sup>

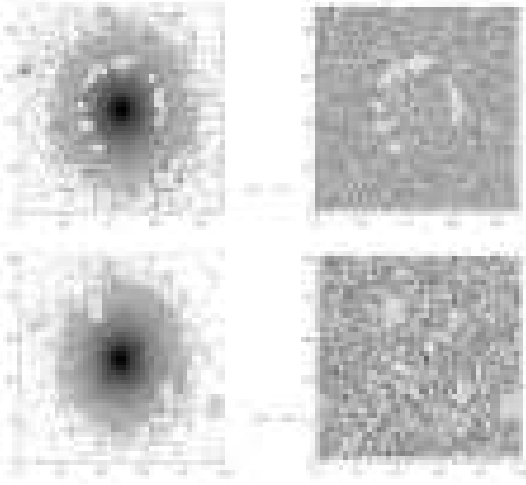
## C2 Morphological classification using concentration index

Galaxies can also be classified according to how ‘peaky’ their light distribution is by using the concentration index (Abraham et al. 1994). The surface brightness distribution of ellipticals and S0s is considerably more centrally concentrated than that for spirals and irregulars. Shimasaku et al. (2001) defined the (inverse) concentration index for SDSS galaxies as the ratio of the half- to the 90% light radii and define an optimum division between late and early types to be at  $C = 0.33$  (with 15 – 20% contamination from opposite types). This separator has also been used to investigate the morphological properties of SDSS galaxies (Goto et al. 2002; Nakamura et al. 2003).

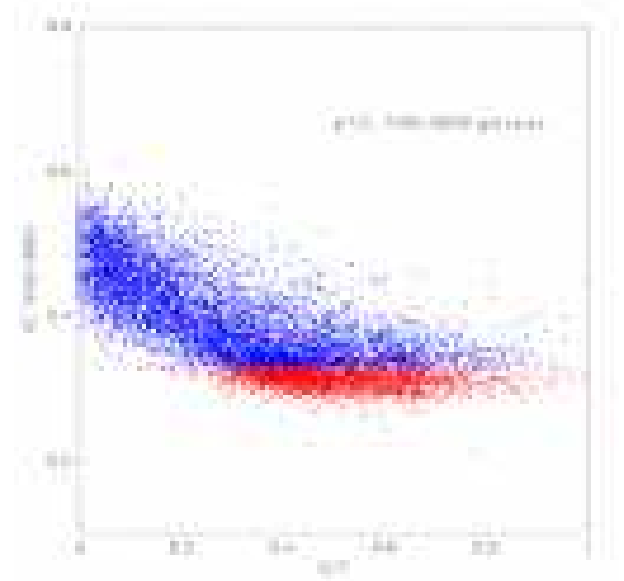


**Figure B3.** Fits to a galaxy that exhibits internal structure. The top images show the unbinned galaxy postage-stamp (left) and the corresponding residual image (right). The bottom images show the galaxy and the residual after the galaxy is binned  $2 \times 2$ . In both cases the  $\chi^2_\nu$  is poor ( $\chi^2_\nu > 2.0$ ) and the residuals are not noise-dominated. This supports the conclusion that galaxies with significant structure give poor fits irrespective of whether they are binned or not.

<sup>5</sup> The red disk population may also include galaxies with heavily obscured star formation.



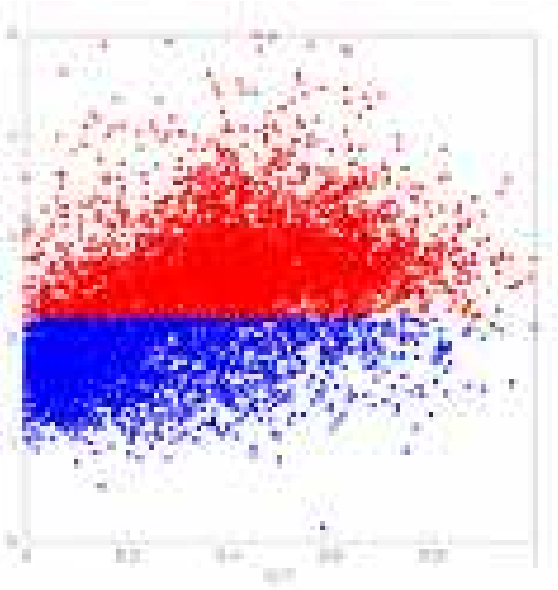
**Figure B4.** Fits to a galaxy that does not exhibit internal structure. The top images show the unbinned galaxy postage-stamp (left) and the corresponding residual image (right). The bottom images show the galaxy and the residual after the galaxy is binned  $2 \times 2$ . In both cases the  $\chi^2_\nu$  is good ( $\chi^2_\nu < 2.0$ ) and the residuals are noise-dominated. This supports the conclusion that galaxies without significant structure result in acceptable fits irrespective of whether they are binned or not.



**Figure C2.** GALACTICA S/T ratio vs. (inverse) concentration index,  $C = R_{50}/R_{90}$ , for our sample of 7493 SDSS galaxies. More centrally concentrated galaxies ( $C < 0.33$ ) predominately have higher S/T ratios.

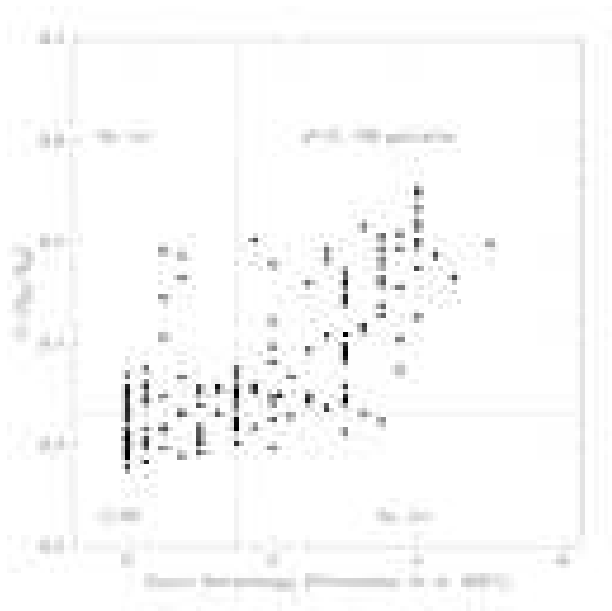
### C3 Morphological classification: S/T vs eye-morphology

Shimasaku et al. (2001) used a sample of 456 bright SDSS galaxies ( $g' < 16.0$ ) visually classified into seven morphological types (Hubble types E, S0, Sa, Sb, Sc, Sdm and Im) to investigate correlations between galaxy colours, effective sizes and concentrations. The (inverse) concentration index was found to correlate well with the visual estimates of morphology. Shimasaku et al. (2001) have kindly provided us with their visual morphologies in order to compare them with the GALACTICA S/T ratios. There are 166 galaxies in common in the two samples which have  $\chi^2_\nu < 2.0$ . Fig. C3 shows that there is a fair correlation between the (inverse) concentration index and the visual morphology for these 166 galaxies, confirming the conclusions of Shimasaku et al. (2001). Plotted in Fig. C4 is the correlation between our derived S/T ratios and the visual morphology for these galaxies. Although the scatter is large, there is a clear trend for the earlier types to have larger S/T ratios.

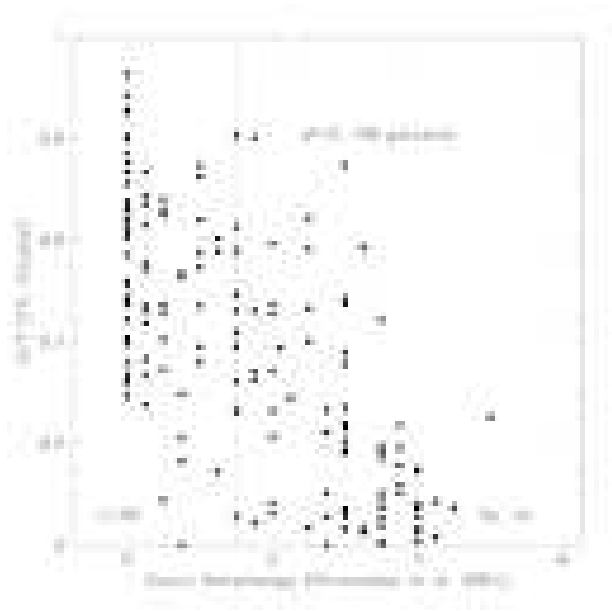


**Figure C1.** GALACTICA S/T ratio vs.  $(u-r)$  colour for the sample of 7493 SDSS galaxies studied here. Most galaxies colour-classified as late types are predominately disk-dominated systems while most galaxies colour-classified as early type have  $S/T > 0.4$ .

Fig. C2 shows that there is a good correlation between the S/T ratios and the (inverse) concentration index,  $C$ . Spheroid-dominated galaxies tend to be more centrally concentrated than disk-dominated galaxies which, however, have a larger scatter in  $C$ .



**Figure C3.** Concentration index vs. eye morphology for 166 galaxies in common between our sample and that of Shimasaku et al. (2001).



**Figure C4.** GALACTICA S/T ratio vs. the visual morphology determined by Shimasaku et al. (2001) for the 166 galaxies in common in the two sample. There is a general trend for the S/T ratio to increase along the S-S0-E morphological sequence but the scatter is large.